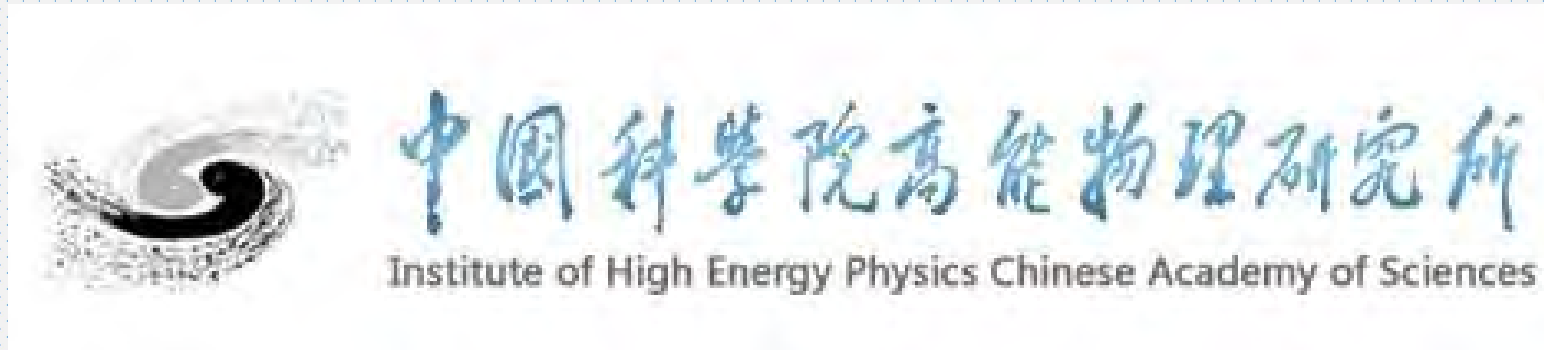


Phase Transitions in an Expanding Universe: Stochastic Gravitational Waves in Standard and Non-Standard Histories

郭怀珂

Aug 11, 2020



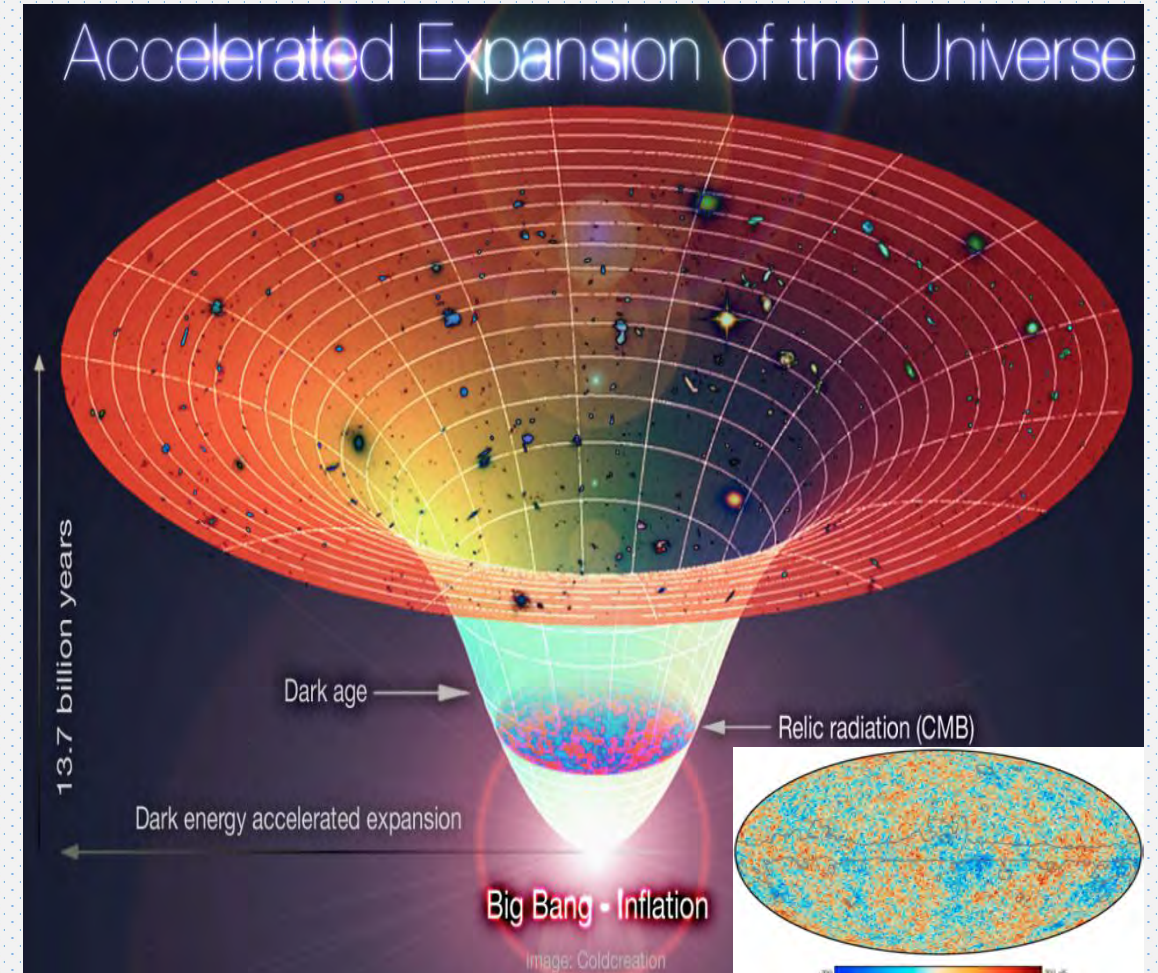
talk based on [arxiv:hep-ph/2007.08537](https://arxiv.org/abs/hep-ph/2007.08537)
in collaboration with Kuver Sinha, Daniel Vagie, Graham White

Triumphs of our Age

Particle Physics

Standard Model of Elementary Particles					
three generations of matter (fermions)			interactions / force carriers (bosons)		
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
QUARKS	u up	c charm	t top	g gluon	H higgs
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
LEPTONS	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

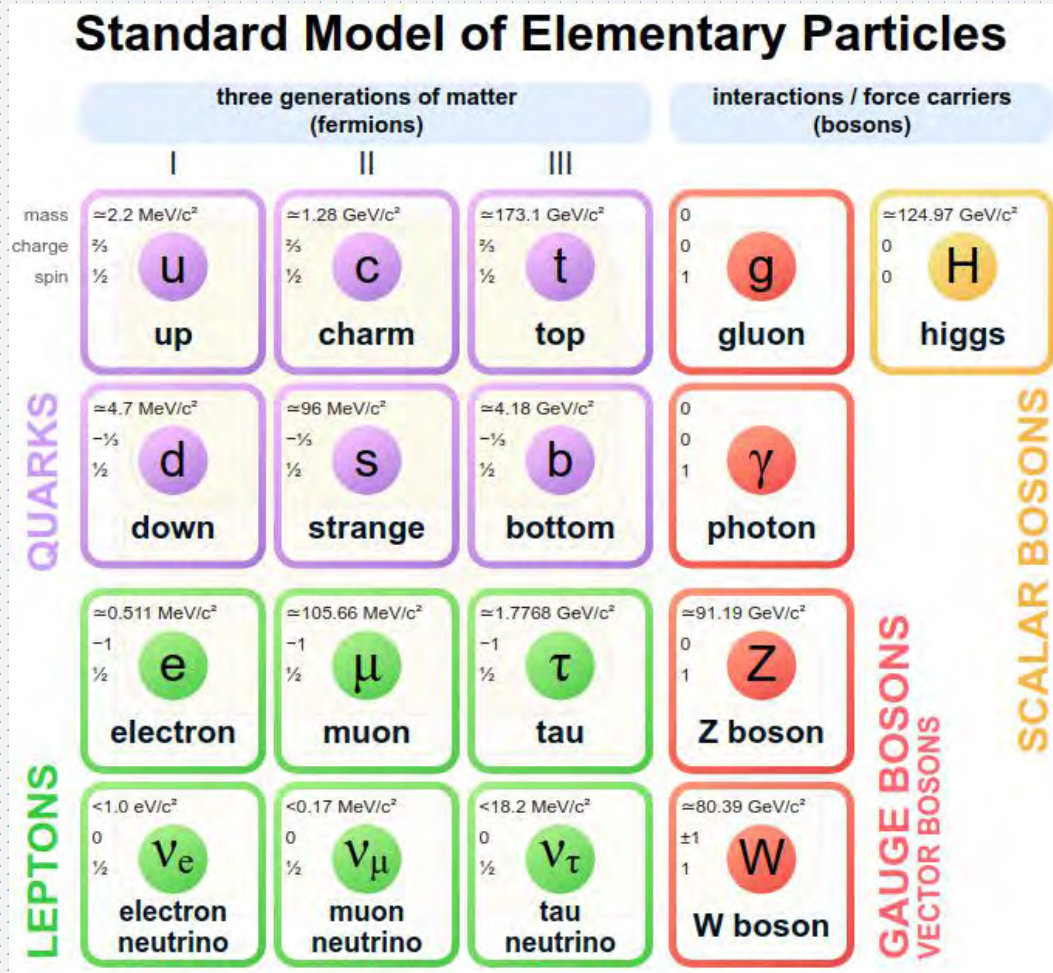
Cosmology



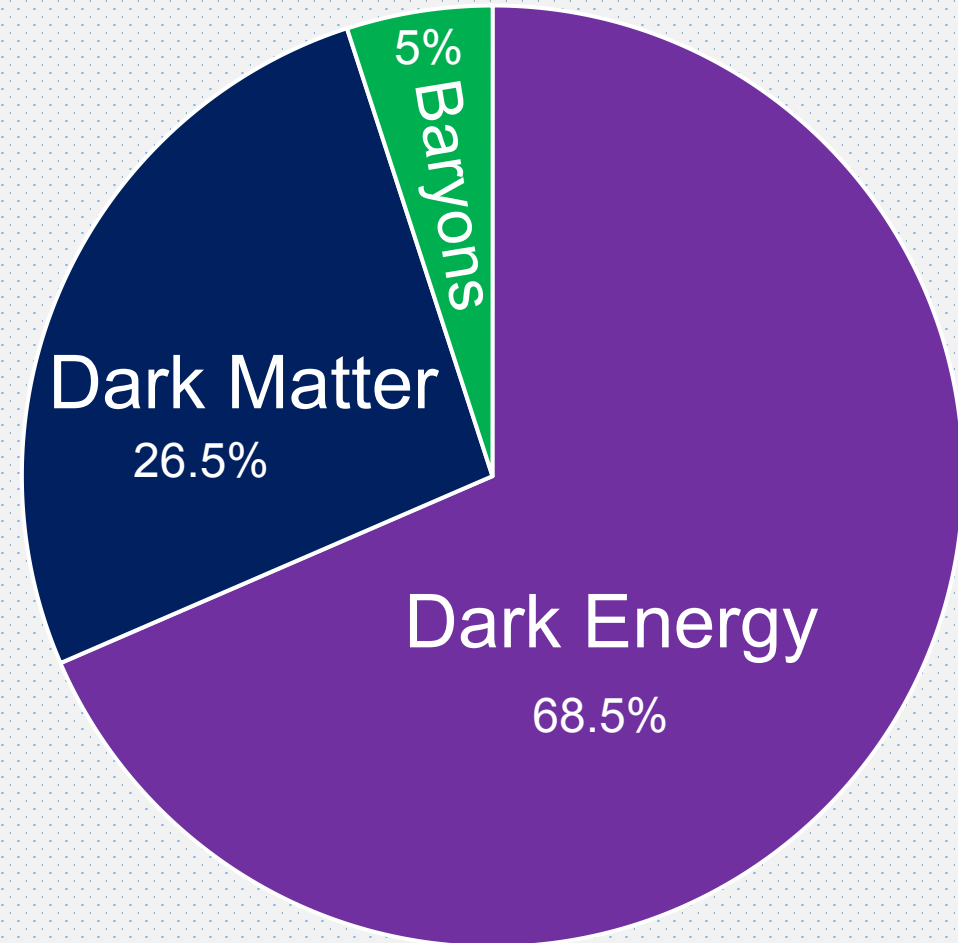
Planck 2018

Triumphs of our Age

Particle Physics



Cosmology

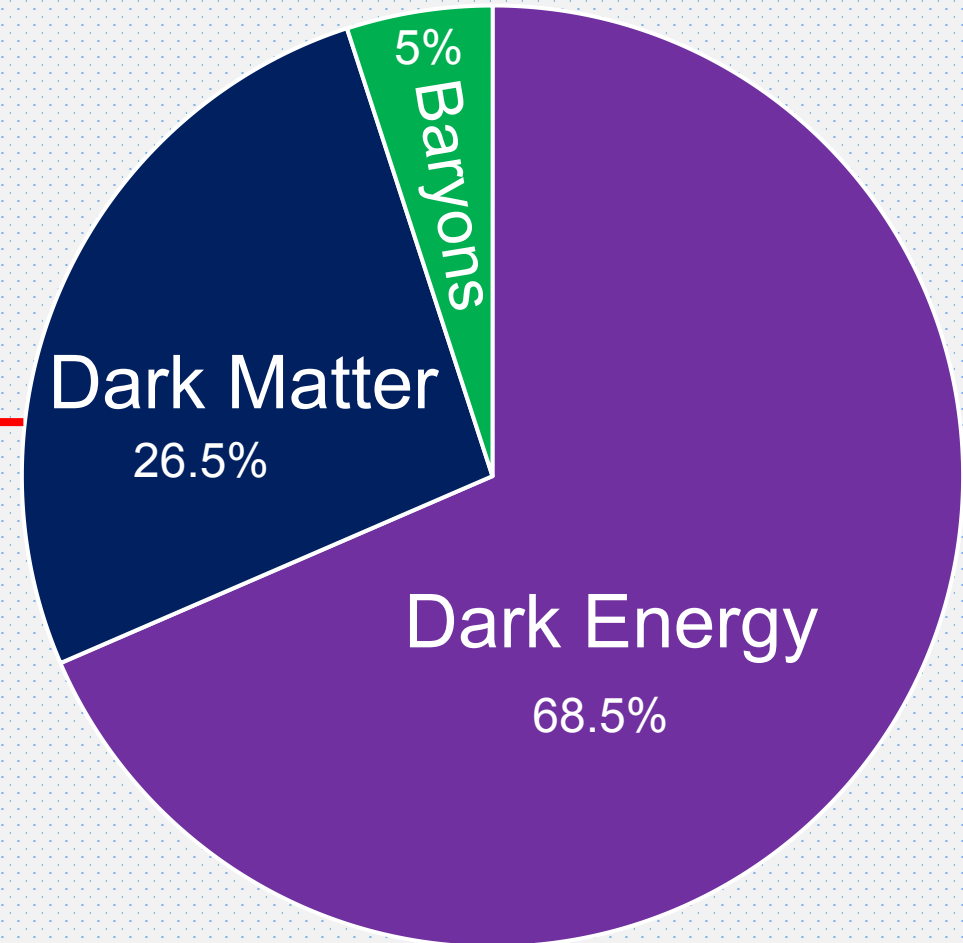


Triumphs of our Age

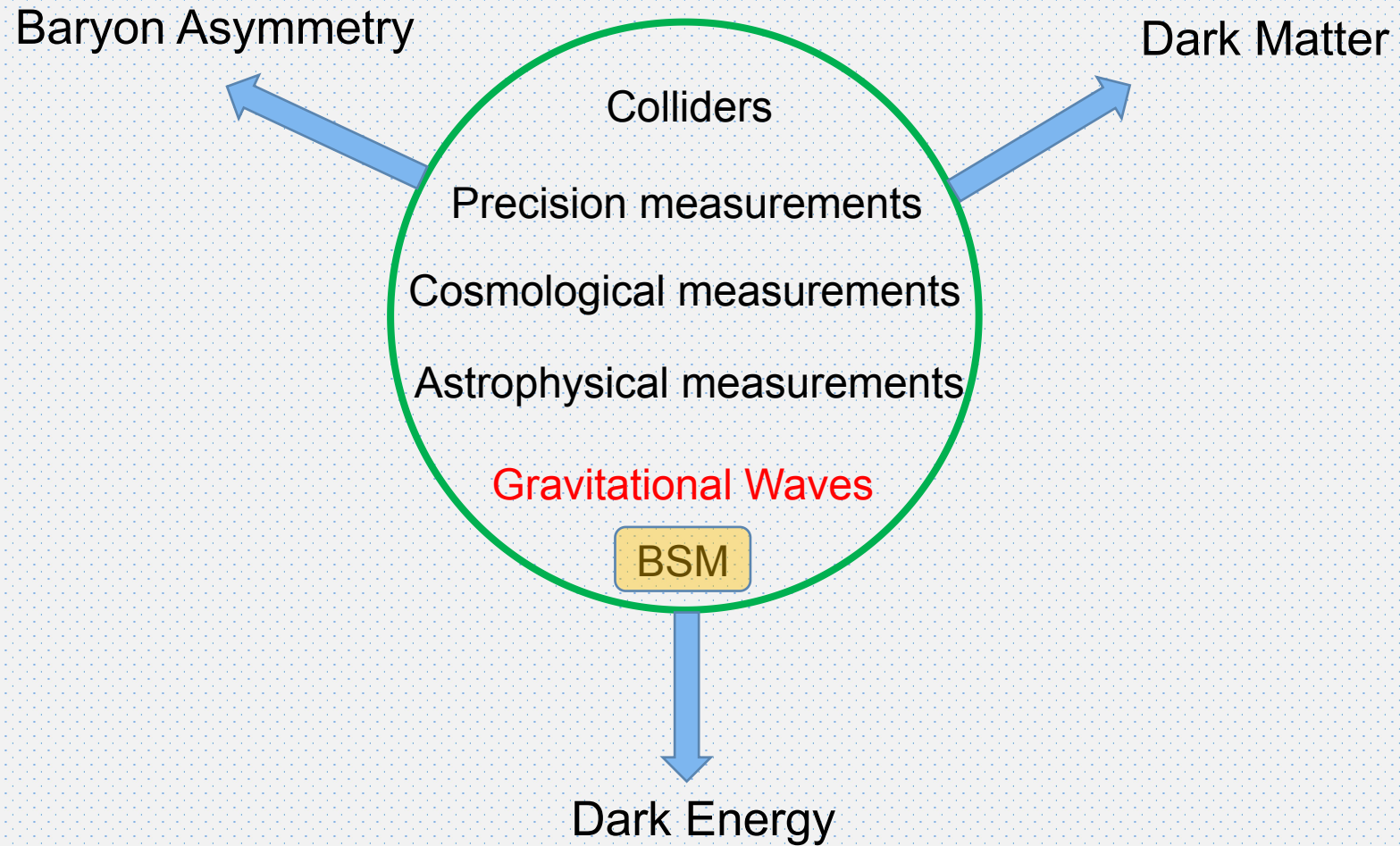
Particle Physics

Standard Model of Elementary Particles					
three generations of matter (fermions)			interactions / force carriers (bosons)		
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
QUARKS	u up	c charm	t top	g gluon	H higgs
	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	d down	s strange	b bottom	γ photon	
LEPTONS	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	e electron	μ muon	τ tau	Z Z boson	
	$< 1.0 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$	
	0	0	0	± 1	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

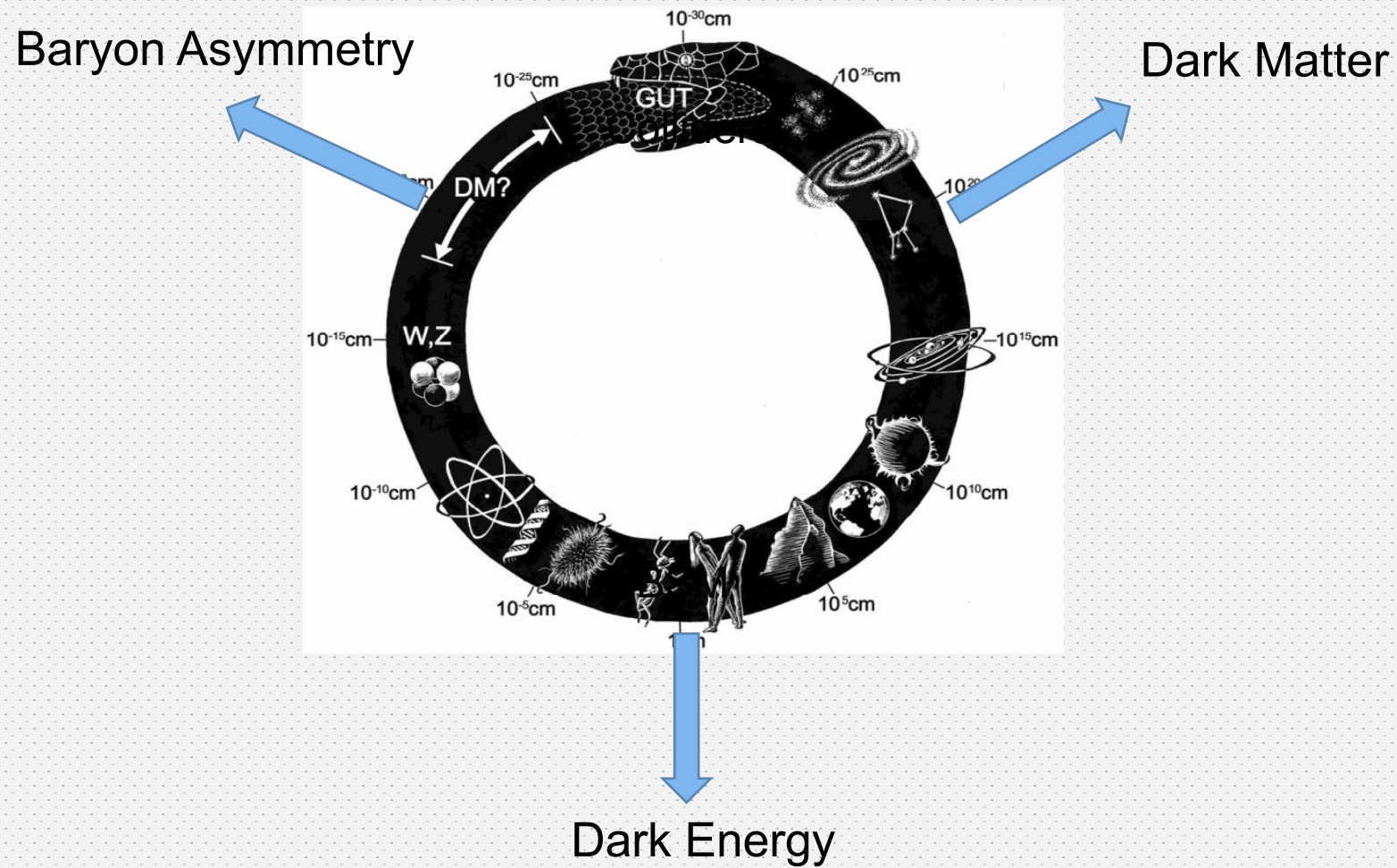
Cosmology



The Problems



The Problems

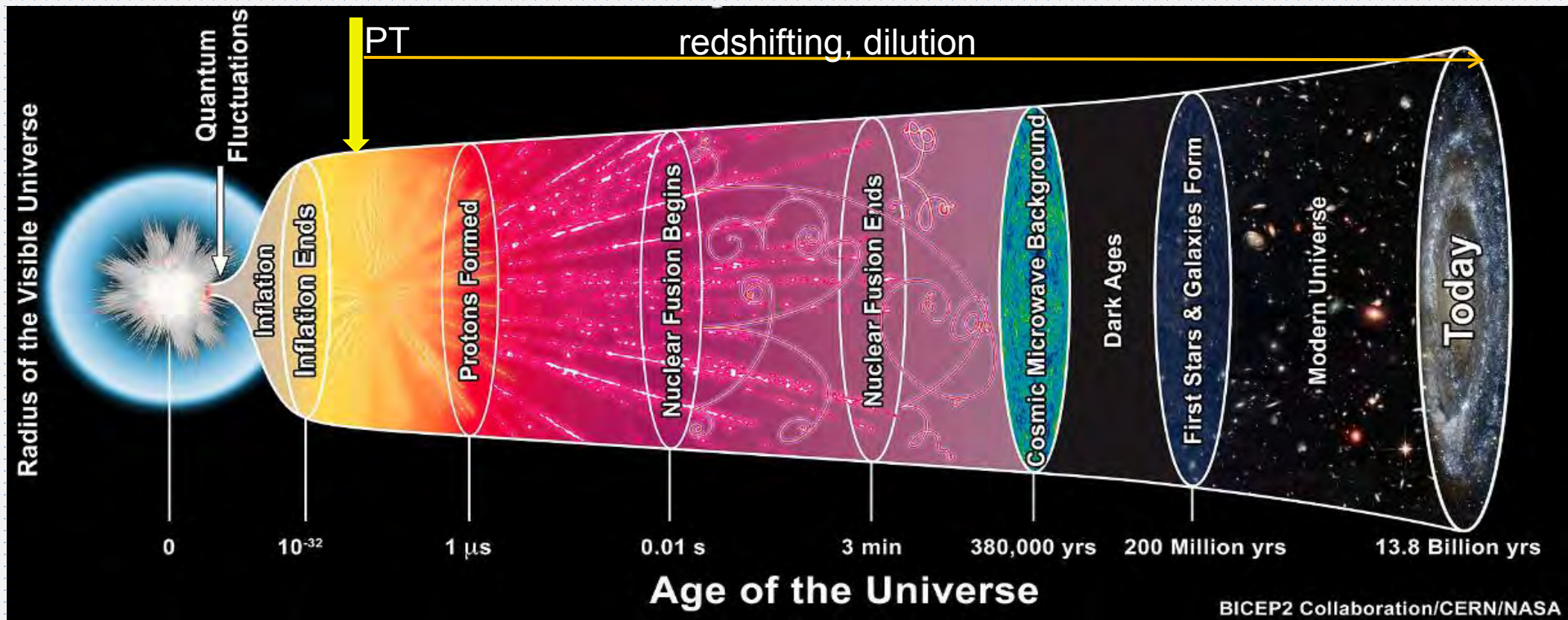


Cosmological Phase Transitions

- Realization of symmetry breakings (Higgs potential, Baryogenesis)
- Gravitational waves as a clean relic, and a cosmic witness

BSM for particle physics: Peccei-Quin, SUSY, Extra Dimensions

BSM for cosmology: Early matter domination(string moduli), Kination, Intermediate Inflationary stage, etc

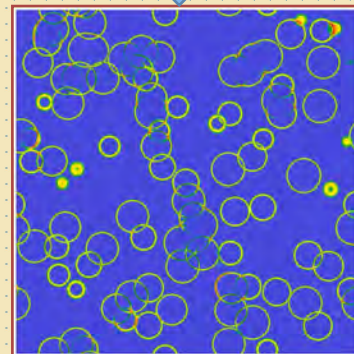
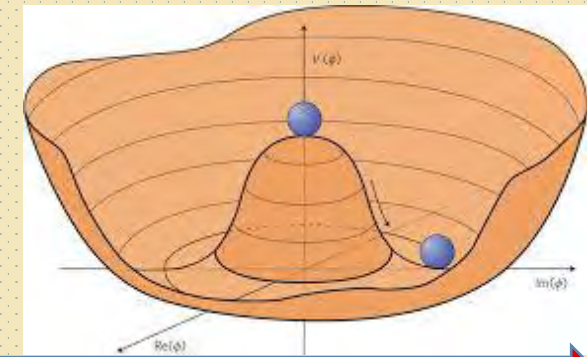
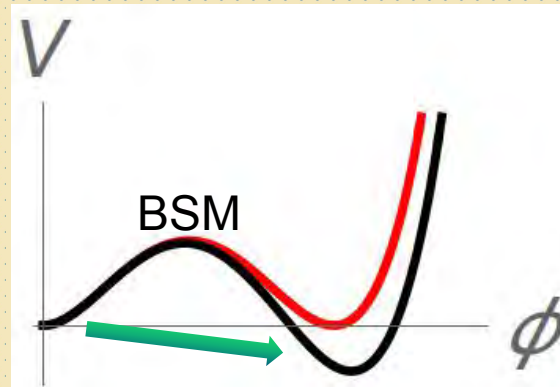
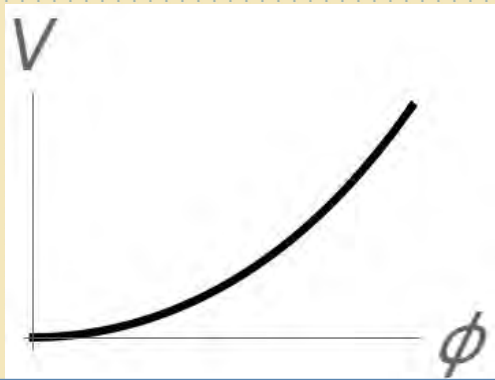


Cosmological Phase Transitions

- Realization of symmetry breakings (Higgs potential, Baryogenesis)
- Gravitational waves as a clean relic, and a cosmic witness

BSM for particle physics: Peccei-Quin, SUSY, Extra Dimensions

BSM for cosmology: Early matter domination(string moduli), Kination, Intermediate Inflationary stage, etc



Hindmarsh, et al, 2015

Cosmological Phase Transitions

- Realization of symmetry breakings (Higgs potential, Baryogenesis)
- Gravitational waves as a clean relic, and a cosmic witness

BSM for particle physics: **Peccei-Quin**, **SUSY**, **Extra Dimensions**

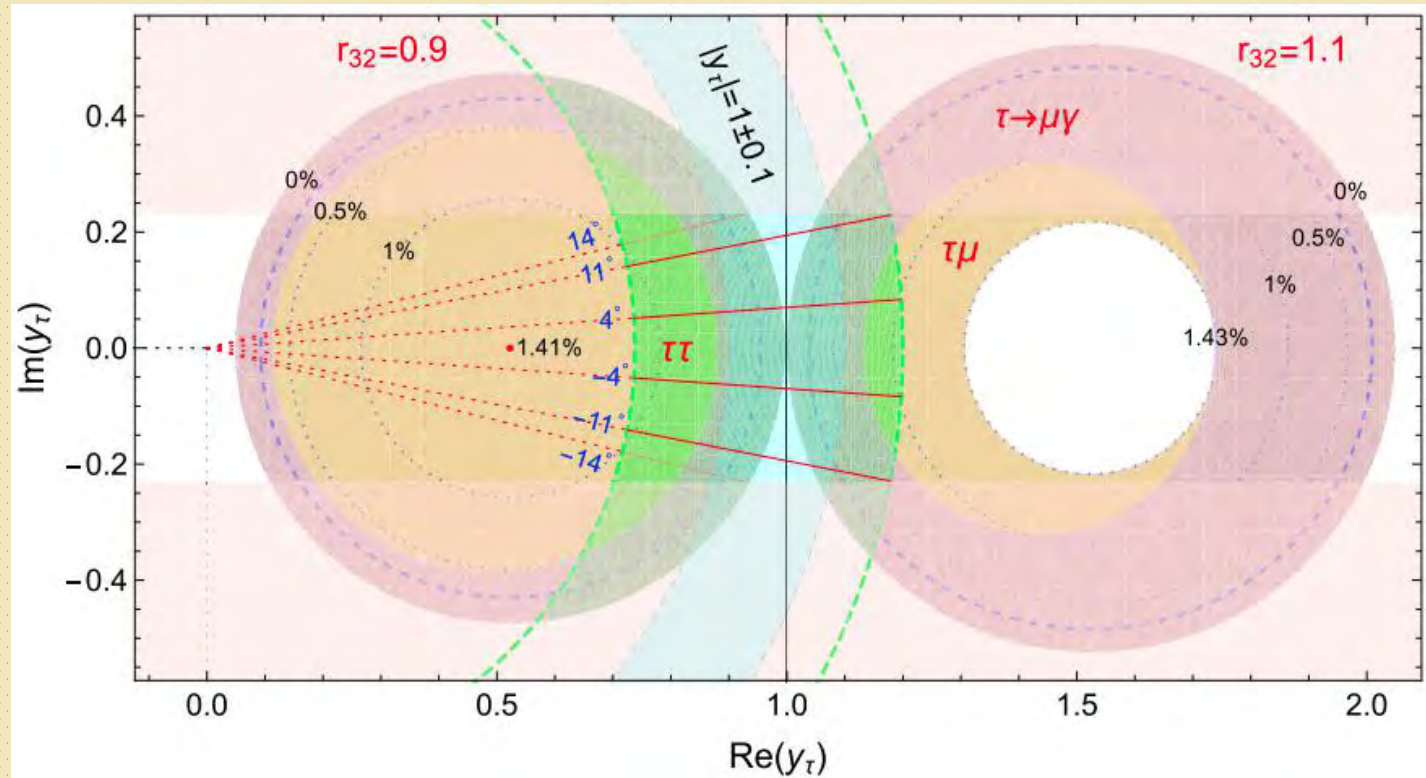
BSM for cosmology: **Early matter domination(string moduli)**, **Kination**, **Intermediate Inflationary stage**, etc

Type III 2HDM

$$\mathcal{L}_{\text{Yukawa}}^{\text{Lepton}} = -\bar{L}^i [Y_{1,ij} \Phi_1 + Y_{2,ij} \Phi_2] e_R^j + \text{H.c.},$$

Jarlskog Invariant

$$J_A = \frac{1}{v^2 \mu_{12}^{\text{HB}}} \sum_{a,b,c=1}^2 v_a v_b^* \mu_{bc} \text{Tr}[Y_c Y_a^\dagger]$$



Cosmological Phase Transitions

- Realization of symmetry breakings (Higgs potential, Baryogenesis)
- Gravitational waves as a clean relic, and a cosmic witness

BSM for particle physics: Peccei-Quin, SUSY, Extra Dimensions

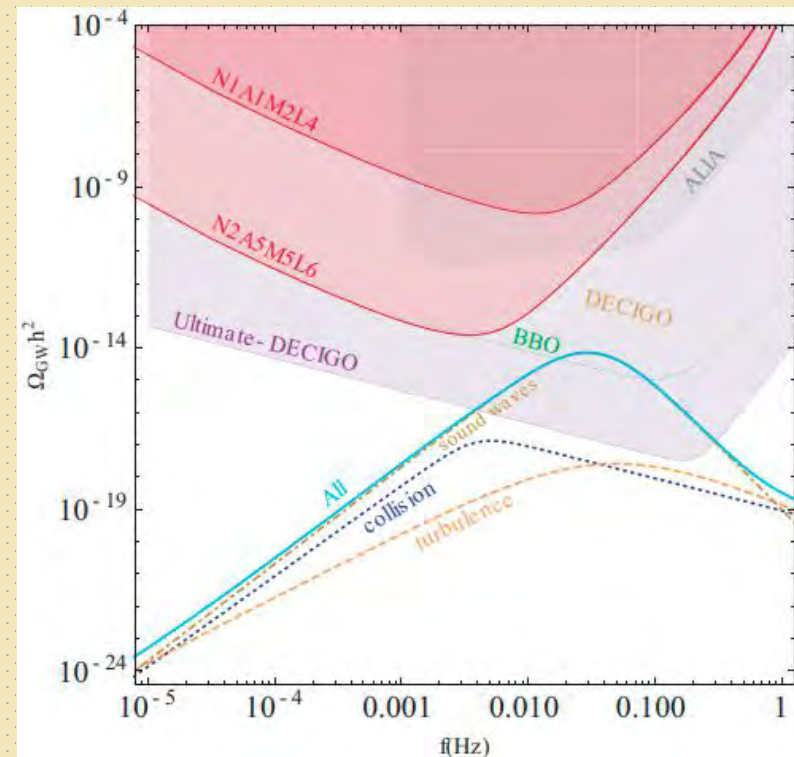
BSM for cosmology: Early matter domination(string moduli), Kination, Intermediate Inflationary stage, etc

NMSSM

$$W_{\text{Higgs}} = \lambda \hat{S} \hat{H}_u \cdot \hat{H}_d + \frac{\kappa}{3} \hat{S}^3,$$
$$-\Delta\mathcal{L}_{\text{soft}} \supset \lambda A_\lambda H_u \cdot H_d S + \frac{1}{3} \kappa A_\kappa S^3 + \text{h.c.}$$

Jarlskog Invariant

$$\mathcal{I} = |\lambda| |\kappa| \sin(\phi'_\lambda - \phi'_\kappa)$$



Bian, Guo, Shu, CPC42,093106(2018)

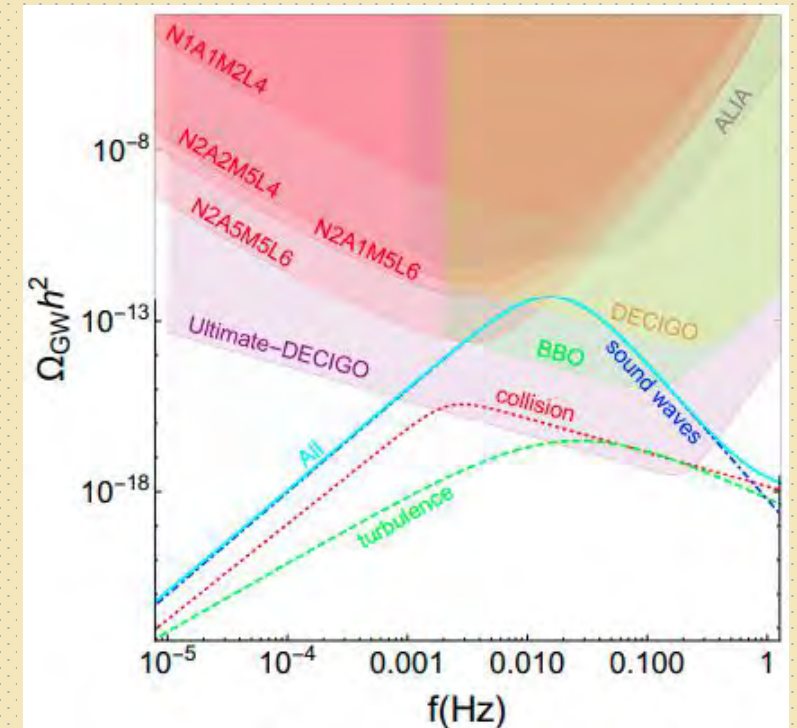
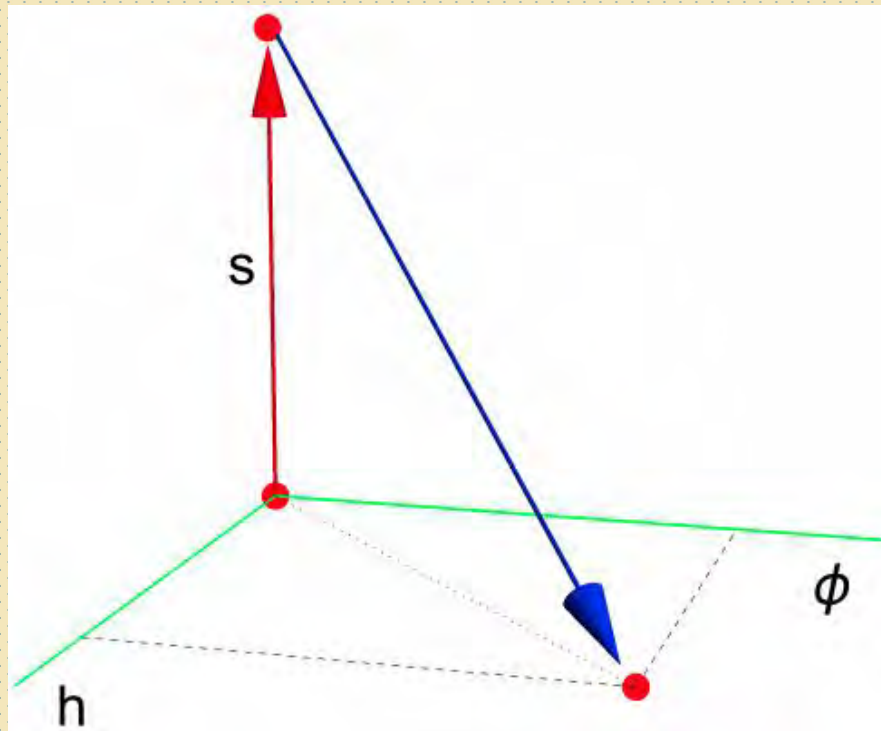
Cosmological Phase Transitions

- Realization of symmetry breakings (Higgs potential, Baryogenesis)
- Gravitational waves as a clean relic, and a cosmic witness

BSM for particle physics: Peccei-Quin, SUSY, Extra Dimensions

BSM for cosmology: Early matter domination(string moduli), Kination, Intermediate Inflationary stage, etc

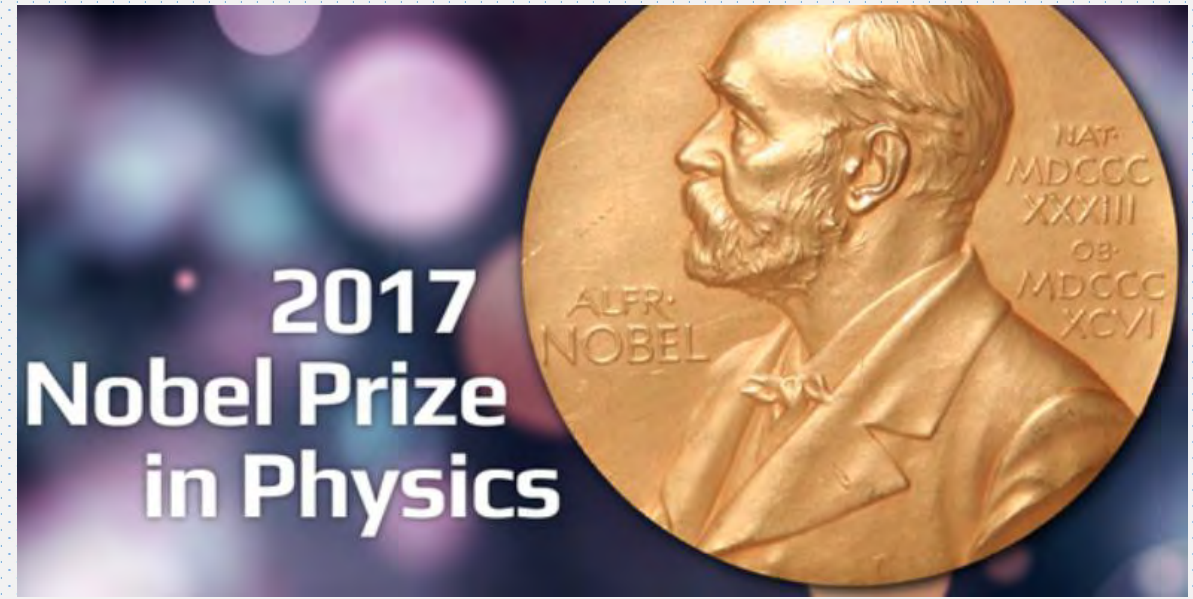
Probe the Dark Matter



LIGO and Gravitational Wave Detection



<https://www.ligo.caltech.edu/>



<https://nobelprize.org/>

A new era for Gravitational Wave Astronomy

Also new tools for particle physics and cosmology

Space-based Gravitational Wave Detectors

Supermassive Black Hole Binary

Extreme Mass Ratio Inspiral

Stochastic Gravitational Waves

arm length: several million-km
mili-hz GW signal



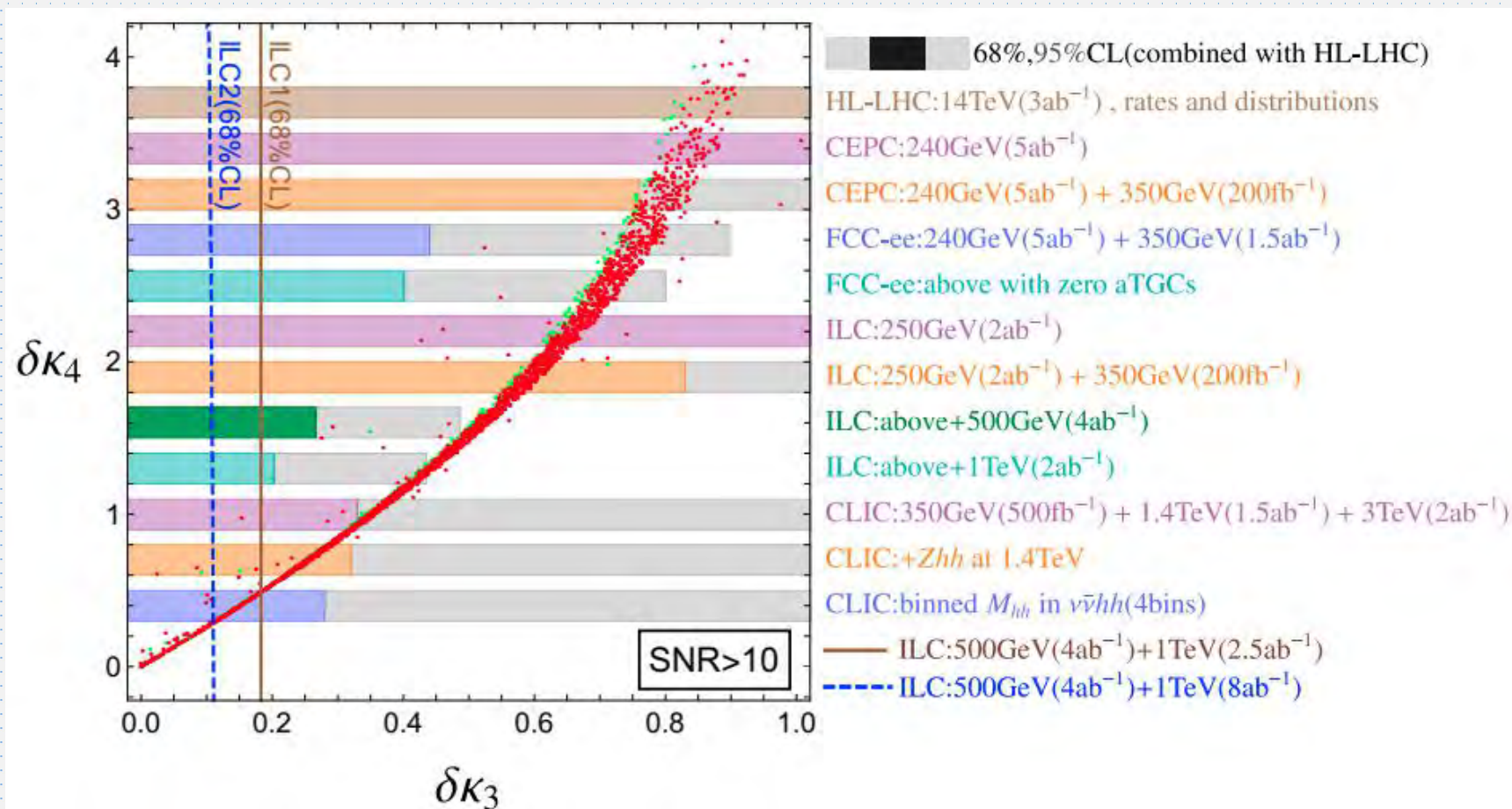
<https://lisa.nasa.gov/>

天琴、太极、LISA、BBO、DECIGO

Collider and GW Correlation and Complementarity

Higgs self-couplings

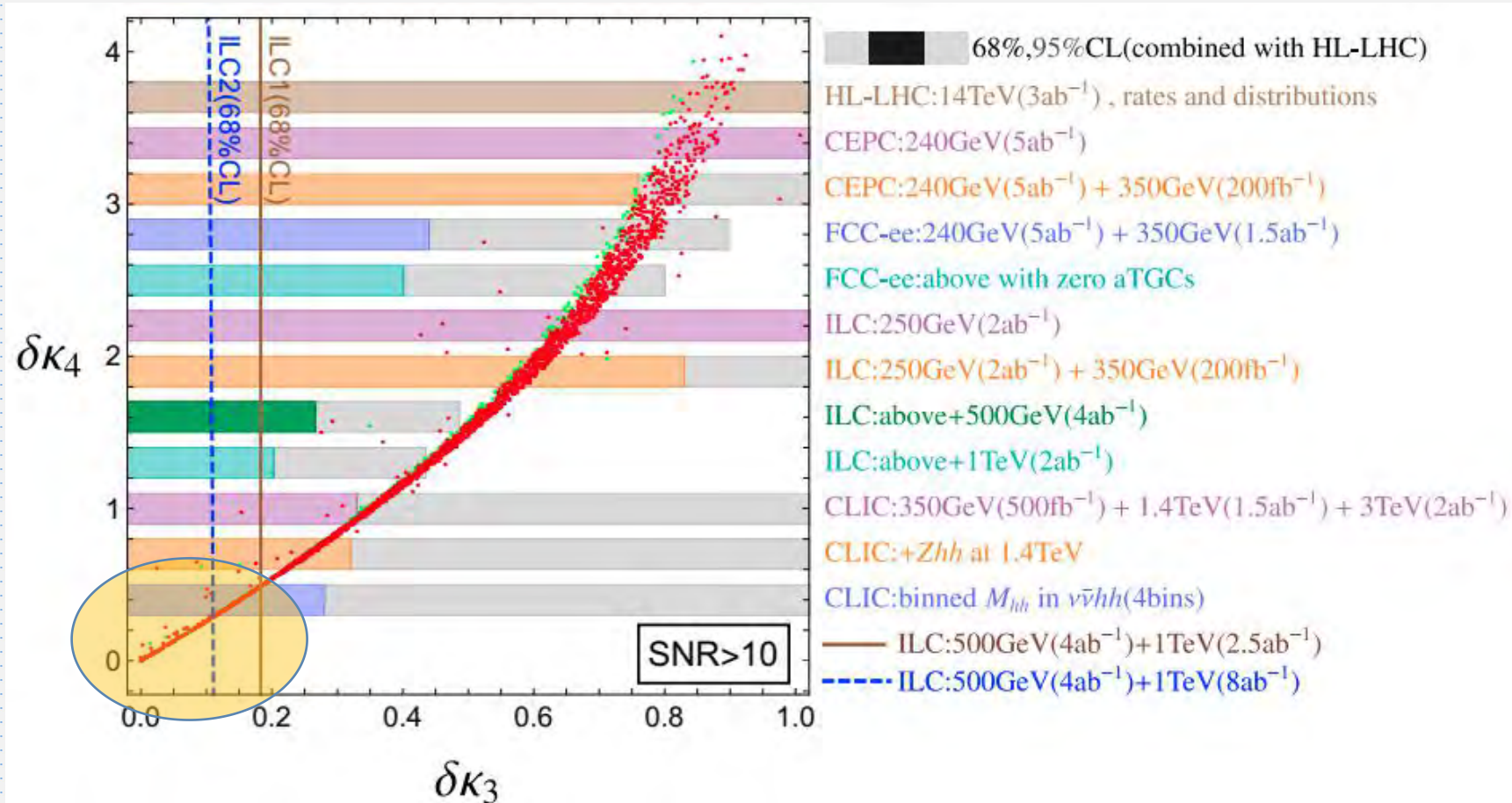
$$\Delta\mathcal{L} = -\frac{1}{2} \frac{m_{h_1}^2}{v} (1 + \delta\kappa_3) h_1^3 - \frac{1}{8} \frac{m_{h_1}^2}{v^2} (1 + \delta\kappa_4) h_1^4$$



Collider and GW Correlation and Complementarity

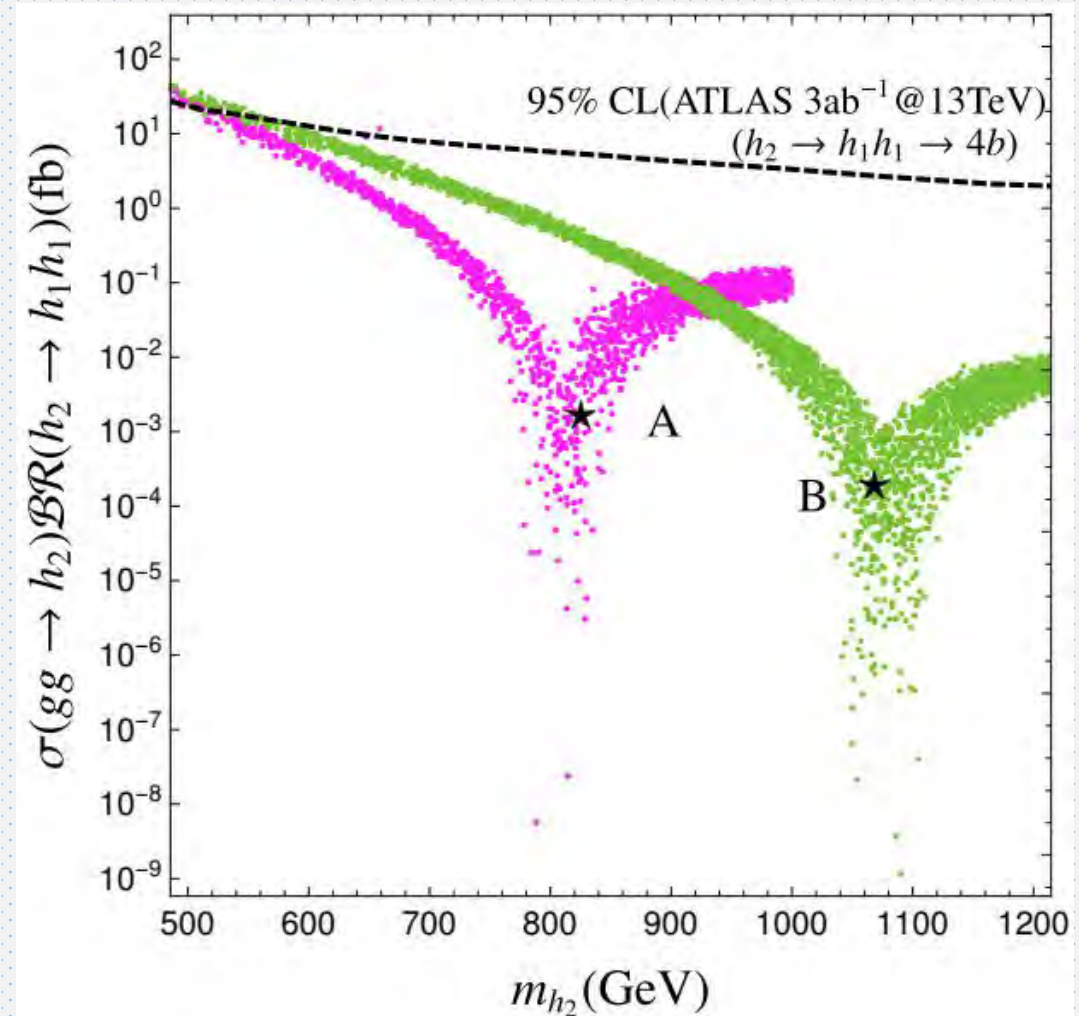
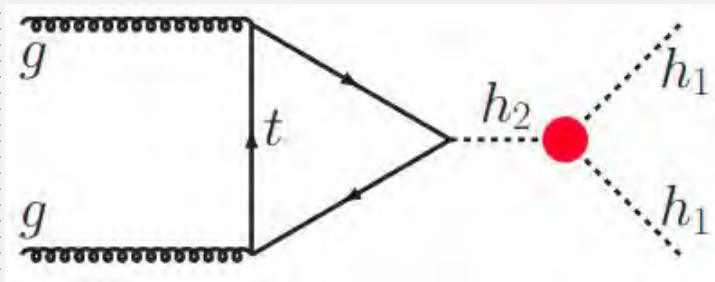
Higgs self-couplings

$$\Delta\mathcal{L} = -\frac{1}{2} \frac{m_{h_1}^2}{v} (1 + \delta\kappa_3) h_1^3 - \frac{1}{8} \frac{m_{h_1}^2}{v^2} (1 + \delta\kappa_4) h_1^4$$



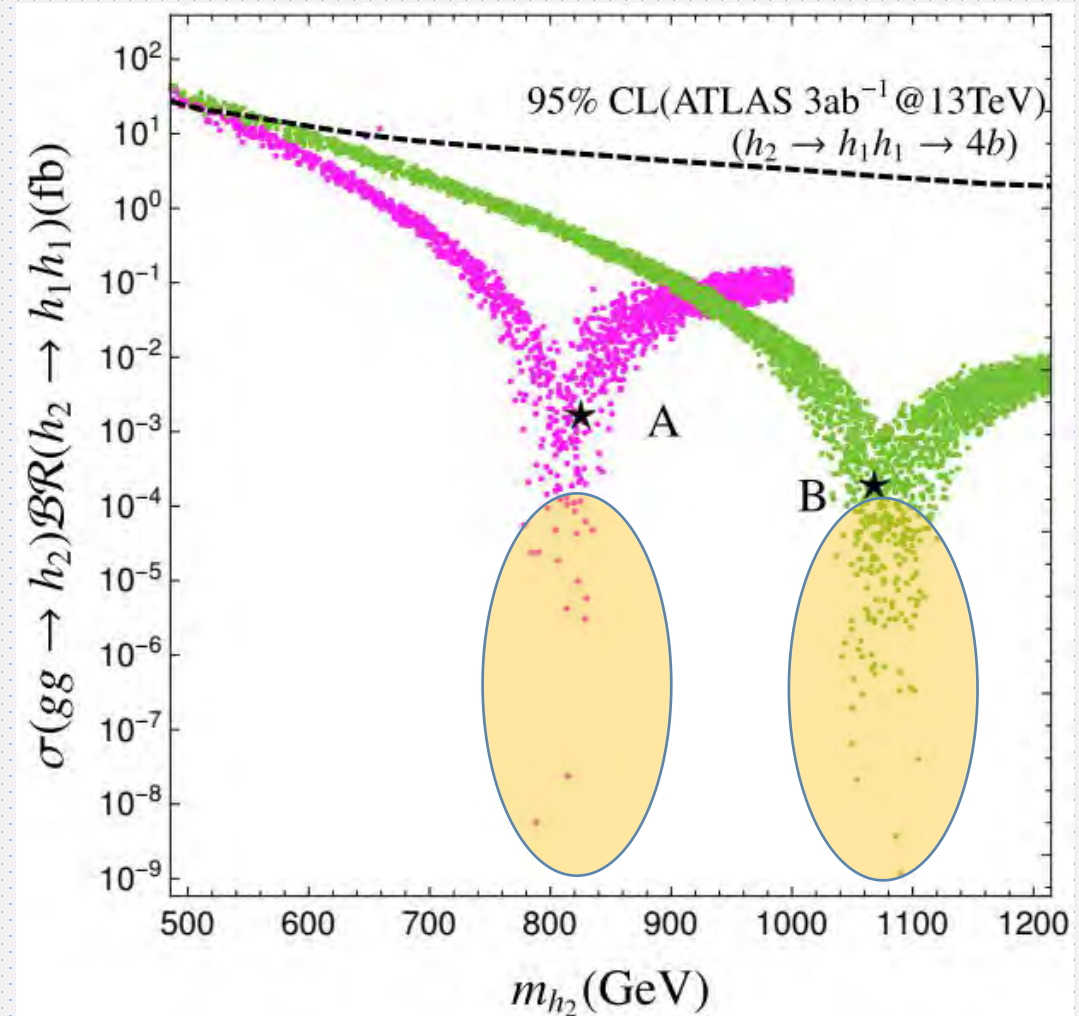
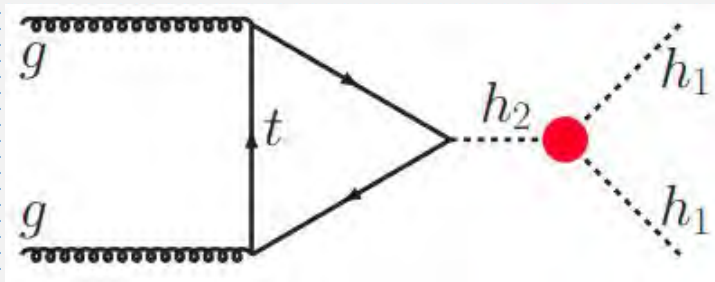
Blind Spots of Collider Measurements

Most striking complementarity
blind spot in Di-Higgs, but strong GWs

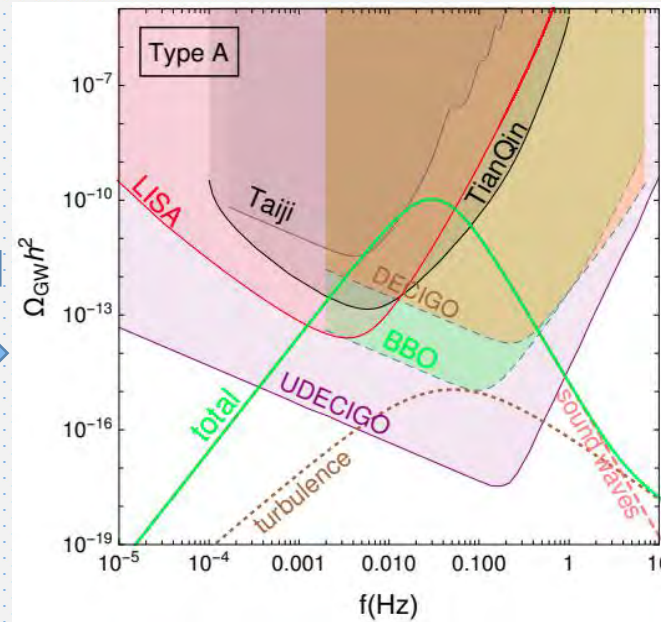
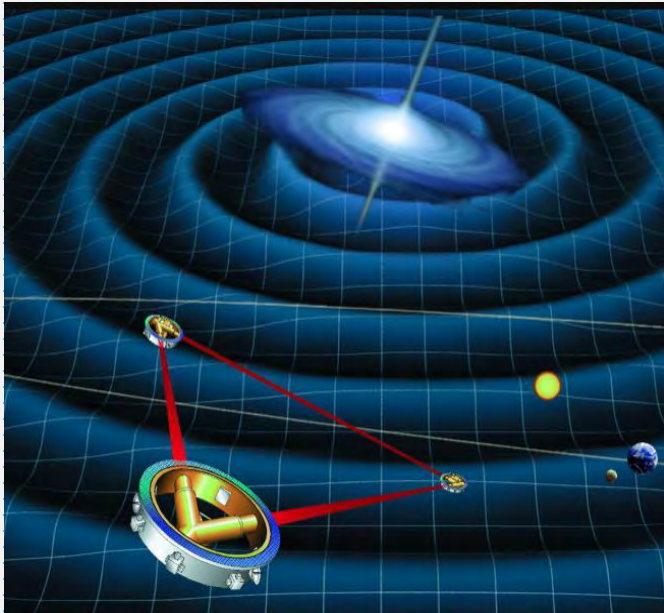


Blind Spots of Collider Measurements

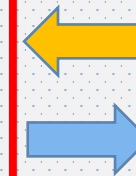
Most striking complementarity
blind spot in Di-Higgs, but strong GWs



The Roadmap



α
 β
 v_w
 T_{end}
 g_s
 \dots

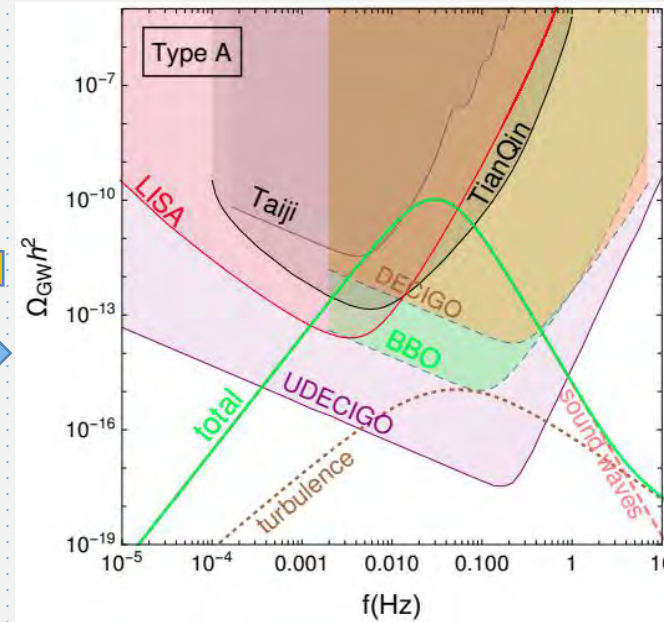
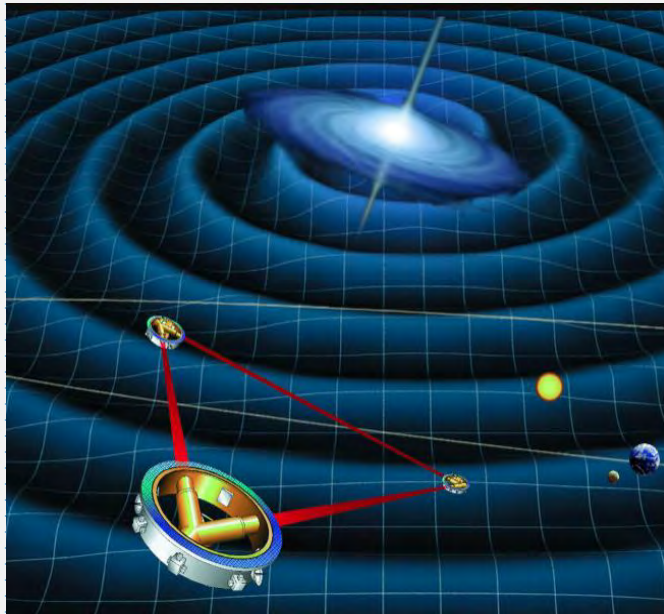


Standard Model of Elementary Particles					
three generations of matter (fermions)			interactions / force carriers (bosons)		
I	II	III			
mass charge spin	$\approx 2.2 \text{ MeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	$\approx 1.28 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	$\approx 173.1 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	0 0 1	$\approx 124.97 \text{ GeV}/c^2$ 0 0
u up	c charm	t top	g gluon		H higgs
$\frac{2}{3}$ $-\frac{1}{3}$ $\frac{1}{2}$	$\frac{2}{3}$ $-\frac{1}{3}$ $\frac{1}{2}$	$\frac{2}{3}$ $-\frac{1}{3}$ $\frac{1}{2}$	0 0 1		
d down	s strange	b bottom	γ photon		BSM
$-\frac{1}{3}$ $-\frac{1}{3}$ $\frac{1}{2}$	$-\frac{1}{3}$ $-\frac{1}{3}$ $\frac{1}{2}$	$-\frac{1}{3}$ $-\frac{1}{3}$ $\frac{1}{2}$			
$\approx 0.511 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	$\approx 105.66 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	$\approx 1.7768 \text{ GeV}/c^2$ -1 $\frac{1}{2}$	$\approx 91.19 \text{ GeV}/c^2$ 0 1		
e electron	μ muon	τ tau	Z Z boson		
$< 1.0 \text{ eV}/c^2$ 0 $\frac{1}{2}$	$< 0.17 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	$< 18.2 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	$\approx 80.39 \text{ GeV}/c^2$ ± 1 1		
ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson		

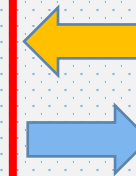
QUARKS
LEPTONS
GAUGE BOSONS
VECTOR BOSONS
SCALAR BOSONS

The Roadmap

theoretical prediction of power spectrum and simulation



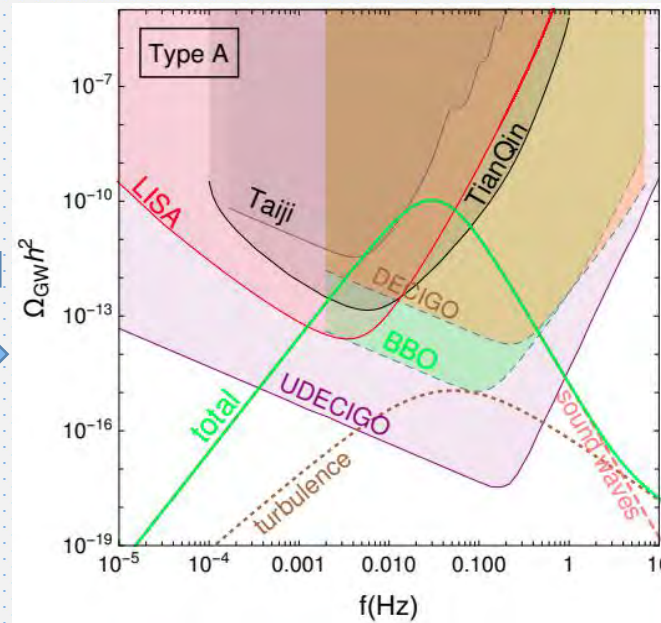
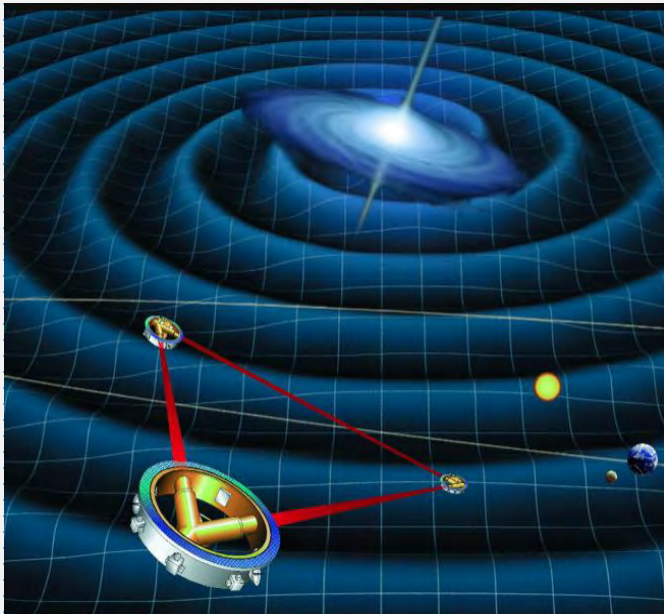
α
 β
 v_w
 T_{end}
 g_s
...



Standard Model of Elementary Particles					
three generations of matter (fermions)			interactions / force carriers (bosons)		
	I	II	III		
mass	≈2.2 MeV/c ²	≈1.28 GeV/c ²	≈173.1 GeV/c ²	0	≈124.97 GeV/c ²
charge	2/3	2/3	2/3	0	0
spin	1/2	1/2	1/2	1	0
	u up	c charm	t top	g gluon	H higgs
	d down	s strange	b bottom	γ photon	BSM
	e electron	μ muon	τ tau	Z Z boson	GAUGE BOSONS
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	VECTOR BOSONS
					SCALAR BOSONS

The Roadmap

theoretical prediction of power spectrum and simulation



α
 β
 v_w
 T_{end}
 g_s
...



Standard Model of Elementary Particles					
three generations of matter (fermions)			interactions / force carriers (bosons)		
I	II	III			
mass $\approx 2.2 \text{ MeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$	mass $\approx 1.28 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$	mass $\approx 173.1 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$	0 0 1	$\approx 124.97 \text{ GeV}/c^2$ 0 0	
u up	c charm	t top	g gluon	H higgs	
mass $\approx 4.7 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$	mass $\approx 96 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$	mass $\approx 4.18 \text{ GeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$	0 0 1		
d down	s strange	b bottom	γ photon		
mass $\approx 0.511 \text{ MeV}/c^2$ charge -1 spin $\frac{1}{2}$	mass $\approx 105.66 \text{ MeV}/c^2$ charge -1 spin $\frac{1}{2}$	mass $\approx 1.7768 \text{ GeV}/c^2$ charge -1 spin $\frac{1}{2}$	0 0 1	$\approx 91.19 \text{ GeV}/c^2$ -1 1	
e electron	μ muon	τ tau	Z Z boson		
mass $< 1.0 \text{ eV}/c^2$ charge 0 spin $\frac{1}{2}$	mass $< 0.17 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$	mass $< 18.2 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$	$\approx 80.39 \text{ GeV}/c^2$ ± 1 1		
ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson		



data analysis, constraints or discovery(parameter estimation)

e.g., LIGO O1, O2 results

Focus of the Following Talk

- Precision calculation of the gravitational wave spectrum

Lay out the framework for modelling GW production in an expanding universe

Any changes to the spectrum?

- Scrutize for hints of non-standard expansion histories

Early matter domination(string moduli), Kination, Intermediate Inflationary stage(supercooling), etc

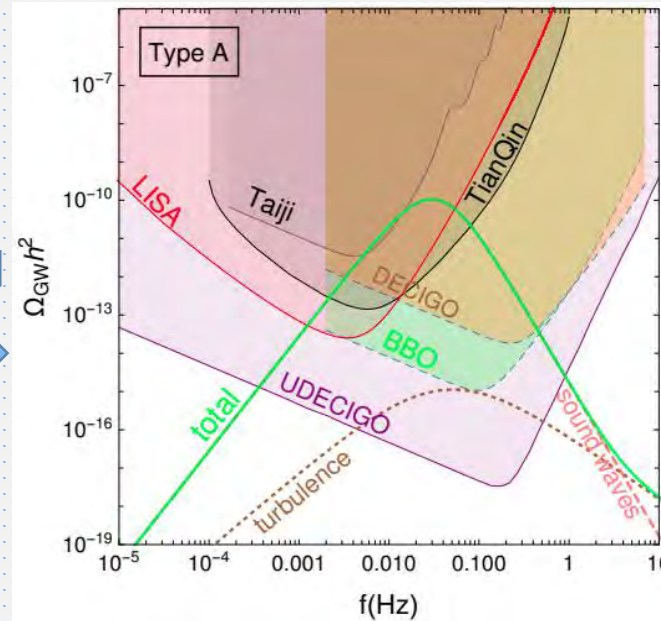
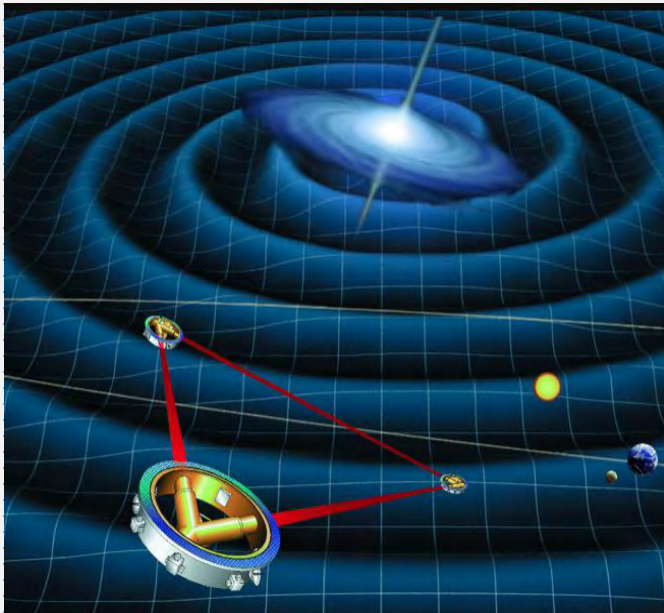
- Do we need a new simulation?

Simulation is costly, and may not be necessary.

Analytical insight of the underlying physics

The Roadmap

theoretical prediction of power spectrum and simulation



α
 β
 v_w
 T_{end}
 g_s
 \dots



Standard Model of Elementary Particles					
three generations of matter (fermions)			interactions / force carriers (bosons)		
I	II	III			
mass charge spin					
$\approx 2.2 \text{ MeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	$\approx 1.28 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	$\approx 173.1 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	0 0 1	0 0 0	$\approx 124.97 \text{ GeV}/c^2$ 0 0
u up	c charm	t top	g gluon		H higgs
$\approx 4.7 \text{ MeV}/c^2$ $-\frac{2}{3}$ $\frac{1}{2}$	$\approx 96 \text{ MeV}/c^2$ $-\frac{2}{3}$ $\frac{1}{2}$	$\approx 4.18 \text{ GeV}/c^2$ $-\frac{2}{3}$ $\frac{1}{2}$	0 0 1		
d down	s strange	b bottom	γ photon		
$\approx 0.511 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	$\approx 105.66 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	$\approx 1.7768 \text{ GeV}/c^2$ -1 $\frac{1}{2}$	0 0 1		
e electron	μ muon	τ tau	Z Z boson		
$< 1.0 \text{ eV}/c^2$ 0 $\frac{1}{2}$	$< 0.17 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	$< 18.2 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	$\approx 80.39 \text{ GeV}/c^2$ ± 1 1		
ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson		

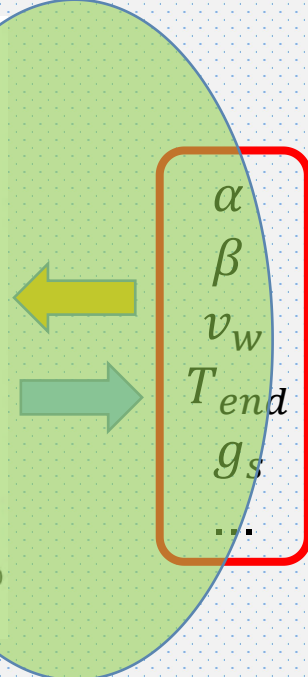
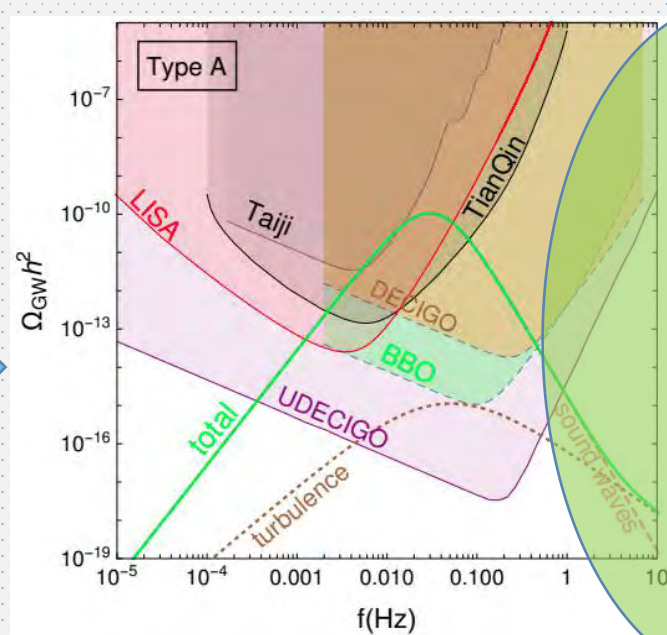
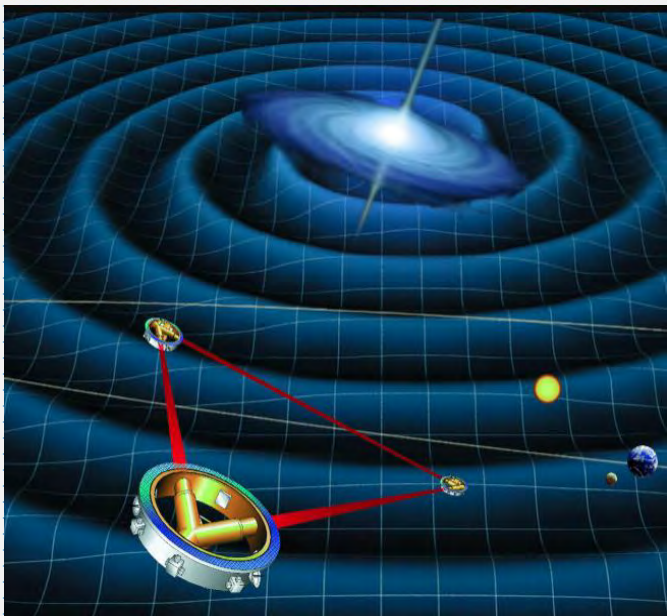


data analysis, constraints or discovery(parameter estimation)

e.g., LIGO O1, O2 results

The Roadmap

theoretical prediction of power spectrum and simulation



Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III		
mass $\approx 2.2 \text{ MeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ u up	mass $\approx 1.28 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ c charm	mass $\approx 173.1 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ t top	0 0 1 g gluon	$\approx 124.97 \text{ GeV}/c^2$ 0 0 0 H higgs
mass $\approx 4.7 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ d down	mass $\approx 96 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ s strange	mass $\approx 4.18 \text{ GeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ b bottom	0 0 1 γ photon	BSM
mass $\approx 0.511 \text{ MeV}/c^2$ charge -1 spin $\frac{1}{2}$ e electron	mass $\approx 105.66 \text{ MeV}/c^2$ charge -1 spin $\frac{1}{2}$ μ muon	mass $\approx 1.7768 \text{ GeV}/c^2$ charge -1 spin $\frac{1}{2}$ τ tau	0 0 1 Z Z boson	
mass $< 1.0 \text{ eV}/c^2$ charge 0 spin $\frac{1}{2}$ ν_e electron neutrino	mass $< 0.17 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$ ν_μ muon neutrino	mass $< 18.2 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$ ν_τ tau neutrino	$\approx 80.39 \text{ GeV}/c^2$ ± 1 1 W W boson	

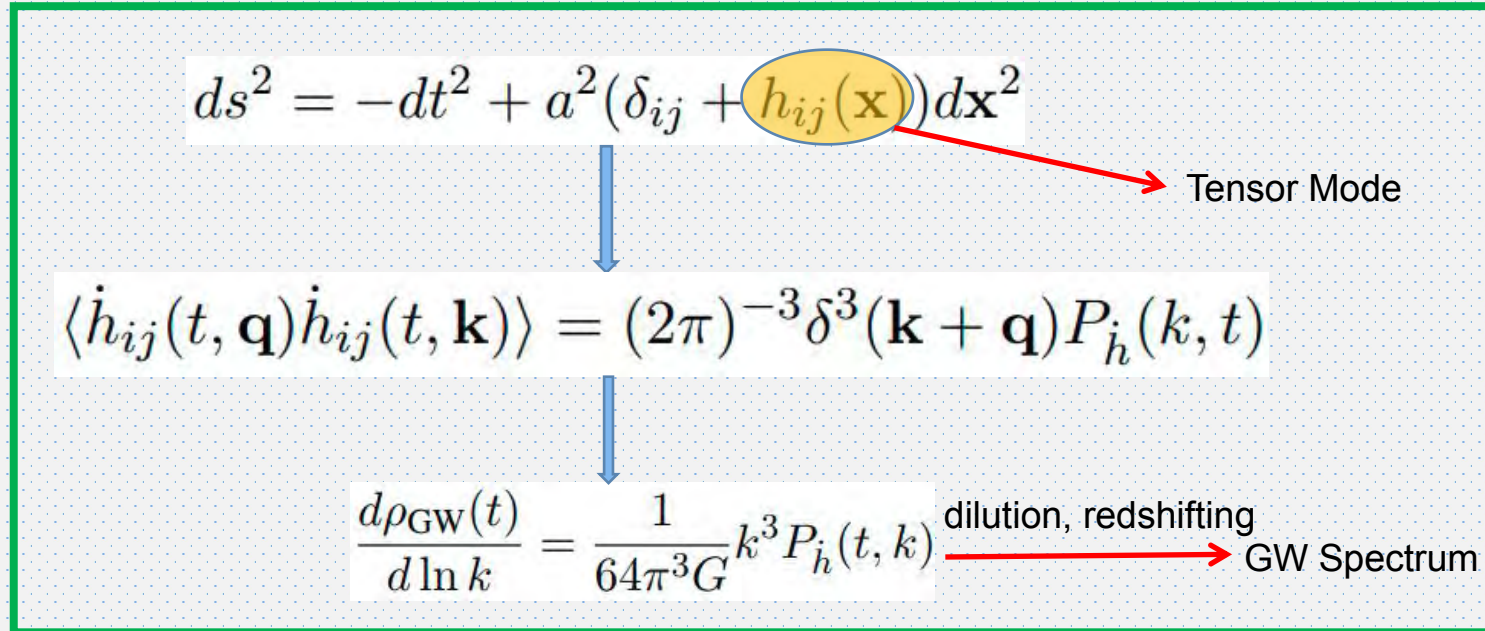
LEPTONS (green box), **QUARKS** (purple box), **GAUGE BOSONS VECTOR BOSONS** (red box), **SCALAR BOSONS** (yellow box).



data analysis, constraints or discovery(parameter estimation)

e.g., LIGO O1, O2 results

How to Calculate Gravitational Waves?



Einstein equation

$$h_q'' + 2\frac{a'}{a}h_q' + q^2 h_q = 16\pi G a^2 \pi_q^T$$

Source evolutions

Plasma(relativistic species), Matter(non-relativistic), Scalar field, EM
 Energy-momentum conservation (hydrodynamic limit)

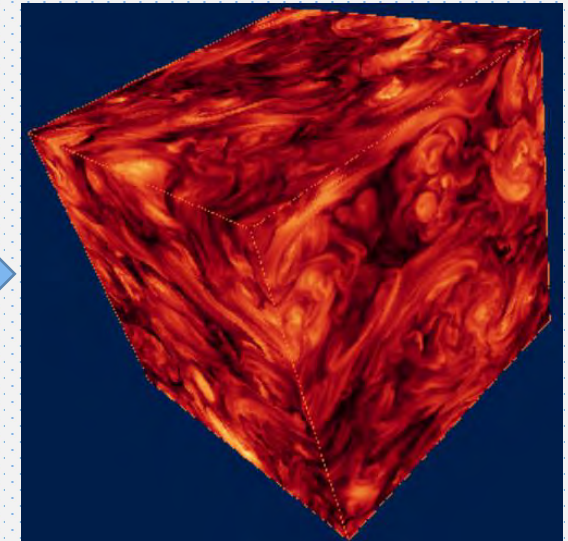
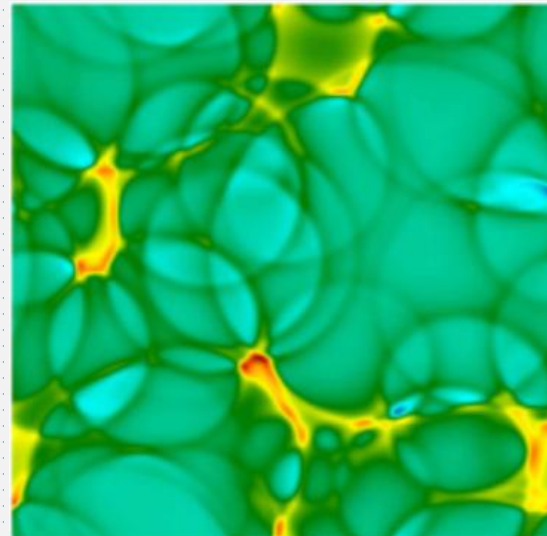
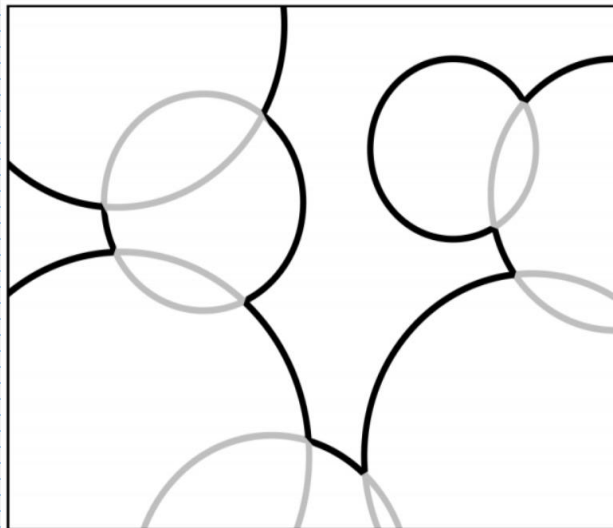
Sources for Gravitational Wave Production

- Bubble Collisions

- Sound Waves in Plasma

dominant in a thermal plasma

- MagnetoHydrodynamic Turbulence



Hindmarsh, et al, PRL112,041301(2013)

<https://home.mpcdf.mpg.de/~wcm/projects/homog-mhd/mhd.html>

Bubble Collisions

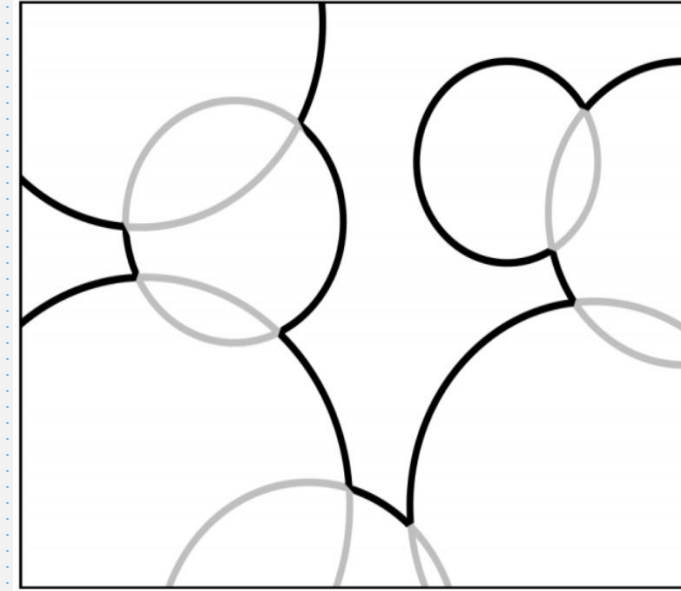
Envelope Approximation

Simulations:

Kosowsky, Turner, Watkins, Kamionkowski
PRL69,2026(1992), PRD45,4514(1992), PRD47,4372(1993), PRD49,2837(1994)
Huber, Konstandin, JCAP09(2008)022

Analytical Modelling:

Jinno, Takimoto, PRD95,024009(2017)



Beyond the Envelope Approximation

Bulk flow model: Konstandin, JCAP03(2018)047, Jinno, Takimoto, JCAP01(2019)060

Direct large scalar lattice simulations: Cutting, Escartin, Hindmarsh, Weir, PRD97,123513(2018), arXiv:2005.13537:

$$h^2 \Omega_{\text{coll}}(f) = 1.67 \times 10^{-5} \Delta(v_w) \left(\frac{H_n}{\beta} \right)^2 \left(\frac{\kappa_\phi \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{1/3} S_{\text{env}}(f)$$

Bubble Collisions

Envelope Approximation

Simulations:

Kosowsky, Turner, Watkins, Kamionkowski
PRL69,2026(1992), PRD45,4514(1992), PRD47,4372(1993), PRD49,2837(1993)
Huber, Konstandin, JCAP09(2008)022

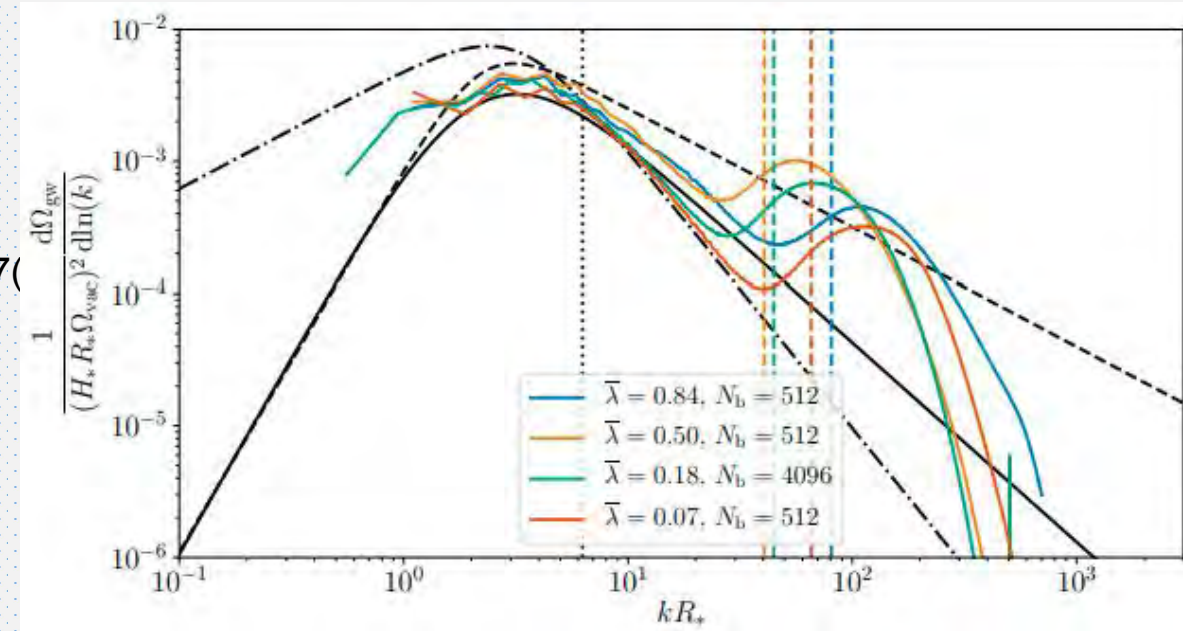
Analytical Modelling:

Jinno, Takimoto, PRD95,024009(2017)

Beyond the Envelope Approximation

Bulk flow model: Konstandin, JCAP03(2018)047, Jinno, Takimoto, JCAP01(2019)060

Direct large scalar lattice simulations: Cutting, Escartin, Hindmarsh, Weir, PRD97,123513(2018), arXiv:2005.13537:



$$h^2 \Omega_{\text{coll}}(f) = 1.67 \times 10^{-5} \Delta(v_w) \left(\frac{H_n}{\beta} \right)^2 \left(\frac{\kappa_\phi \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{1/3} S_{\text{env}}(f)$$

Magnetohydrodynamic Turbulence

Analytical Modelling

Kolmogorov spectrum:

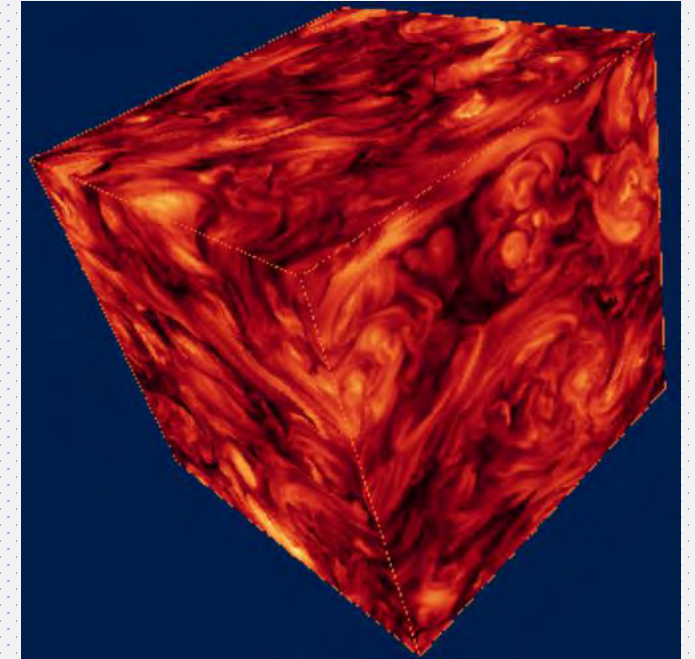
Kosowsky, Mack, Kahniashvili, PRD66,024030(2002)

Gogoberidze, Kahniashvili, Kosowsky, PRD76,083002(2007)

Caprini, Durrer, Servant, JCAP12(2009)024

Numerical Simulations

Pol, Mandal, Brandenburg, Kahniashvili, Kosowsky, arxiv:1903.08585



<https://home.mpcdf.mpg.de/~wcm/projects/homog-mhd/mhd.html>

unknown

$$h^2 \Omega_{\text{turb}}(f) = 3.35 \times 10^{-4} \left(\frac{H_*}{\beta} \right) \left(\frac{\kappa_{\text{turb}} \alpha}{1 + \alpha} \right)^{\frac{3}{2}} \left(\frac{100}{g_*} \right)^{1/3} v_w S_{\text{turb}}(f)$$

Sound Waves

Numerical Simulations:

Hindmarsh, Huber, Rummukainen, Weir,
PRL112, 041301 (2014), PRD92, 123009 (2015), PRD96, 103520 (2017)
Reduction found: Cutting, Hindmarsh, Weir, PRL125, 021302 (2020)

Analytical Modelling(sound shell model)

Hindmarsh, 120, 071301 (2018)
Hindmarsh, Hijazi, JCAP12(2019)062

$$h^2 \Omega_{\text{sw}}(f) = 2.65 \times 10^{-6} \left(\frac{H_n}{\beta} \right) \left(\frac{\kappa_v \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{\frac{1}{3}} \times v_w S_{\text{sw}}(f)$$

LISA Cosmology Workinggroup, JCAP04(2016)001

The dominant source in a thermal plasma.

Sound Waves

Numerical Simulations:

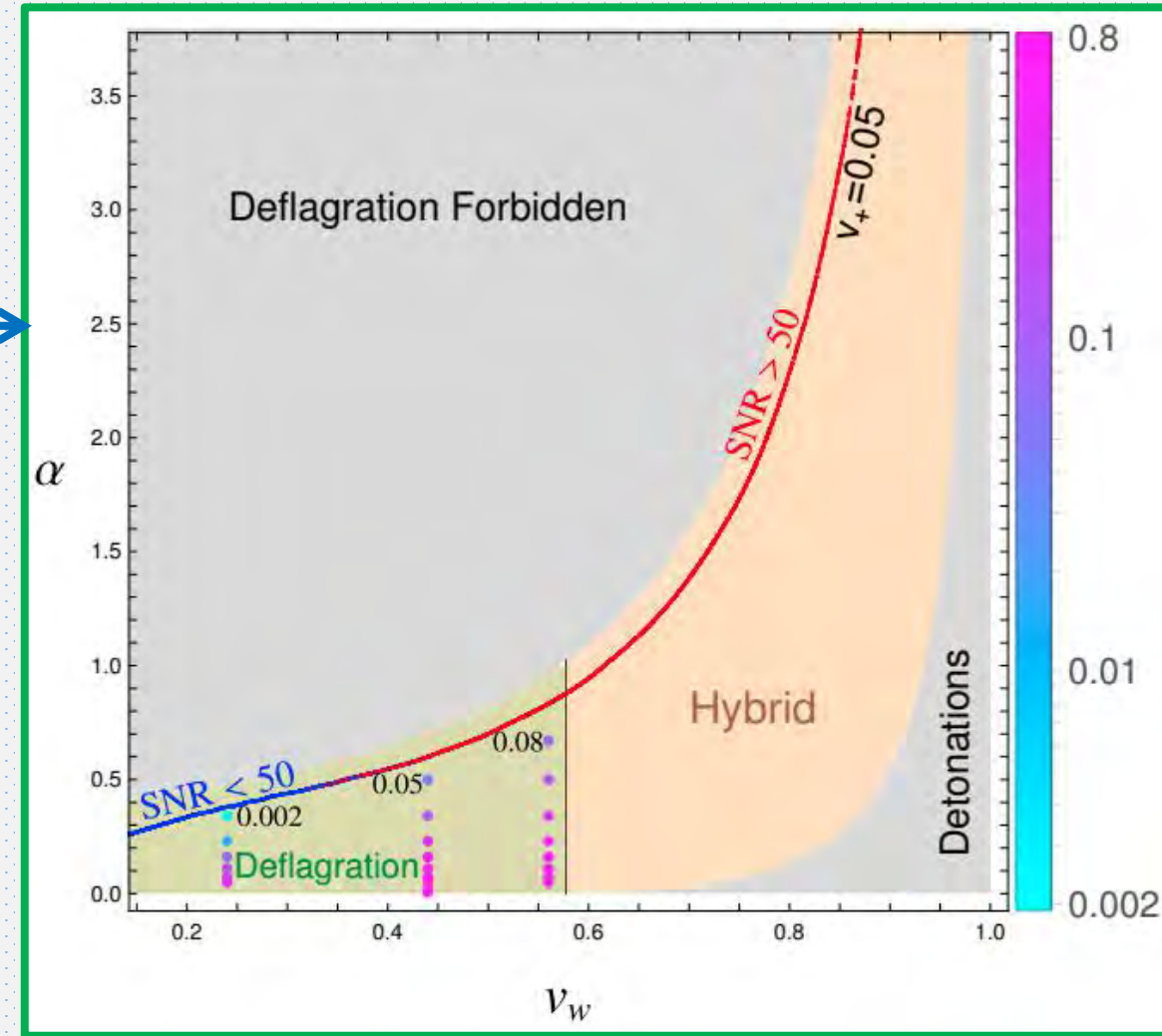
Hindmarsh, Huber, Rummukainen, Weir,
PRL112, 041301 (2014), PRD92, 123009 (2015), PRD96, 103520 (2017)

Reduction found: Cutting, Hindmarsh, Weir, PRL125, 021302 (2020)

Analytical Modelling(sound shell model)

Hindmarsh, 120, 071301 (2018)

Hindmarsh, Hijazi, JCAP12(2019)062



Alves, Goncalves, Ghost, **Guo**, Sinha, JHEP03(2020)053

Sound Waves

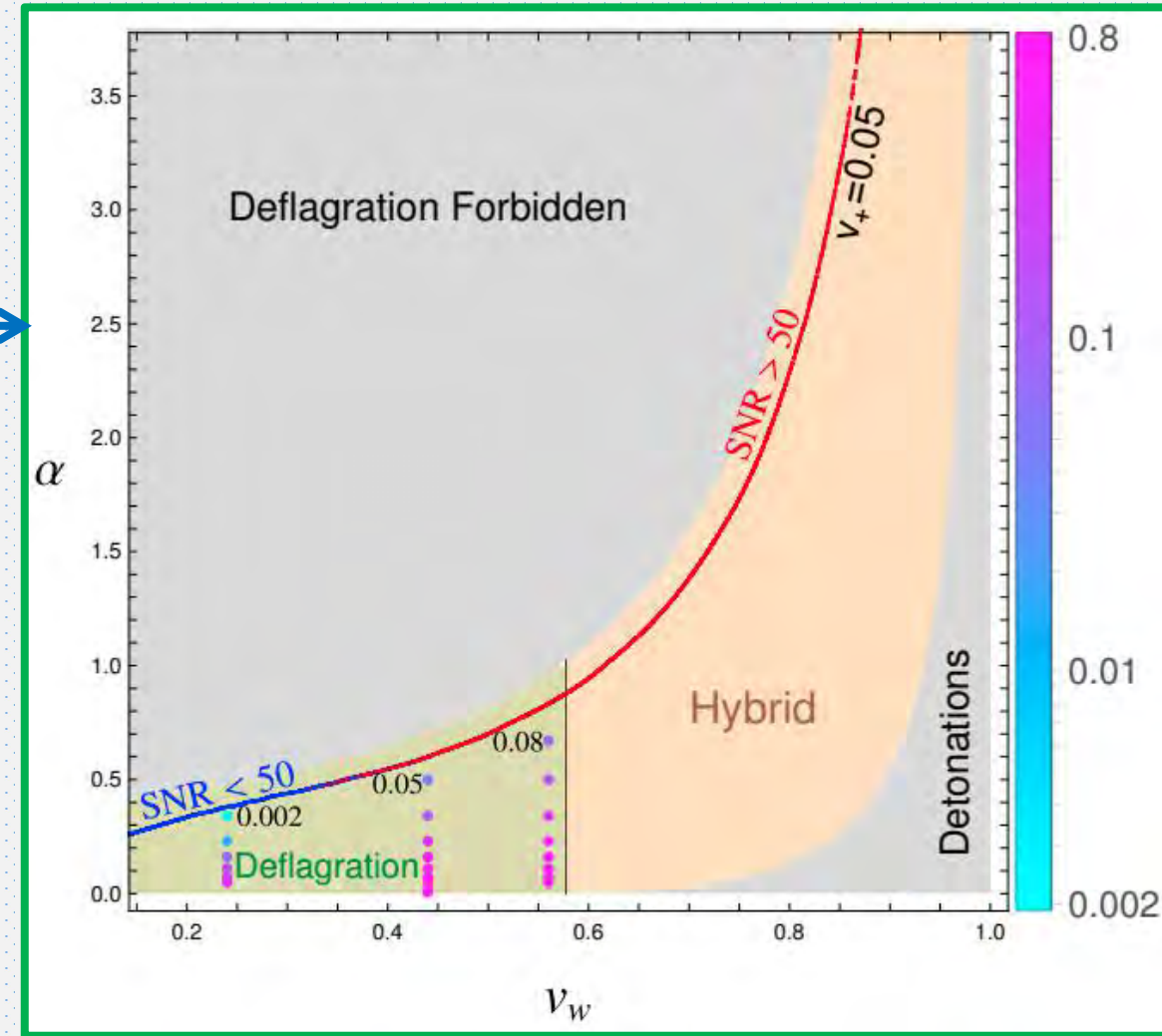
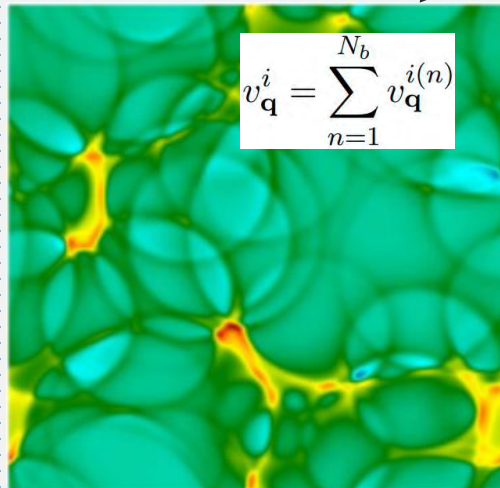
Numerical Simulations:

Hindmarsh, Huber, Rummukainen, Weir,
PRL112, 041301 (2014), PRD92, 123009 (2015), PRD96, 103520 (2017)

Reduction found: Cutting, Hindmarsh, Weir, PRL125, 021302 (2020)

Analytical Modelling(sound shell model)

Hindmarsh, 120, 071301 (2018)
Hindmarsh, Hijazi, JCAP12(2019)062



Alves, Goncalves, Ghost, **Guo**, Sinha, JHEP03(2020)053

Sound Waves

Numerical Simulations:

Hindmarsh, Huber, Rummukainen, Weir,
PRL112, 041301 (2014), PRD92, 123009 (2015), PRD96, 103520 (2017)
Reduction found: Cutting, Hindmarsh, Weir, PRL125, 021302 (2020)

Analytical Modelling(sound shell model)

Hindmarsh, 120, 071301 (2018)
Hindmarsh, Hijazi, JCAP12(2019)062

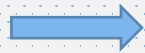
Expanding Universe Analysis(Guo,Sinha,Vagie,White,arxiv:2007.08537):
Numerical simulations: equations in an expanding universe
Analytical modelling in an expanding universe(sound shell model)
Found an additional effect not captured in previous spectrum

How numerical simulations are performed?

- Realized only several years ago (Hindmarsh, et al, PRL 112, 041301, 2013)

$$T^{\mu\nu}_{;\mu}|_{\text{field}} = (\partial^2 \phi) \partial^\nu \phi + \frac{1}{\sqrt{g}} (\partial_\mu \sqrt{g}) (\partial^\mu \phi) (\partial^\nu \phi) - \frac{\partial V}{\partial \phi} \partial^\nu \phi = \delta^\nu,$$
$$T^{\mu\nu}_{;\mu}|_{\text{fluid}} = \partial_\mu [(e + p) U^\mu U^\nu] + \left[\frac{1}{\sqrt{g}} (\partial_\mu \sqrt{g}) g^\nu_\lambda + \Gamma^\nu_{\mu\lambda} \right] (e + p) U^\mu U^\lambda + g^{\mu\nu} \partial_\mu p + \frac{\partial V}{\partial \phi} \partial^\nu \phi = -\delta^\nu.$$

bubble generation



bubble, fluid, metric evolution



gravitational wave measurement



partial differential equation solving on a lattice, difficult

Equations of Motion in an Expanding Universe

Scalar field:

(Klein-Gordon equation but with friction)

$$-\ddot{\phi} + \frac{1}{a^2} \nabla^2 \phi - \frac{\partial V}{\partial \phi} - 3 \frac{\dot{a}}{a} \dot{\phi} = \eta \gamma \left(\dot{\phi} + \frac{1}{a} \mathbf{v} \cdot \nabla \phi \right)$$

Parallel projection:

(Energy equation)

$$\begin{aligned} \dot{E} + p \left[\dot{\gamma} + \frac{1}{a} \nabla \cdot (\gamma \mathbf{v}) \right] + \frac{1}{a} \nabla \cdot (E \mathbf{v}) - \gamma \frac{\partial V}{\partial \phi} \left(\dot{\phi} + \frac{1}{a} \mathbf{v} \cdot \nabla \phi \right) + 3 \frac{\dot{a}}{a} \gamma (e + p) \\ = \eta \gamma^2 \left(\dot{\phi} + \frac{1}{a} \mathbf{v} \cdot \nabla \phi \right)^2, \end{aligned}$$

$$E \equiv e \gamma$$

Vector equation:

(Euler equation)

$$\dot{Z}^i + \frac{1}{a} \nabla \cdot (\mathbf{v} Z^i) + 5 \frac{\dot{a}}{a} Z^i + \frac{1}{a^2} \partial_i p + \frac{1}{a^2} \frac{\partial V}{\partial \phi} \partial_i \phi = - \frac{1}{a^2} \eta \gamma \left(\dot{\phi} + \frac{1}{a} \mathbf{v} \cdot \nabla \phi \right) \partial_i \phi,$$

$$Z^i \equiv \gamma (e + p) U^i = \gamma^2 (e + p) v^i / a.$$

Can not simply reduce to Minkowski form.

But collision phase is short and the expansion effect should be small. **New numerical simulations not needed.** 34

Sound Waves when bubbles have all disappeared

- Equations of motion can be obtained by simply rescaling of Minkowski counterpart
- Sound waves(fluctuations of energy, pressure, velocity)

$$\begin{aligned}(a^4 S^i)' + \nabla \cdot (a^4 S^i \mathbf{v}) + \partial_i(a^4 p) &= 0, & S^i &= \gamma^2(\epsilon + p)v^i \\(a^4 \epsilon \gamma)' + [\gamma' + \nabla \cdot (\gamma \mathbf{v})](a^4 p) + \nabla \cdot (a^4 \epsilon \gamma \mathbf{v}) &= 0, \\ \gamma^2(v' + \frac{1}{2} \hat{\mathbf{v}} \cdot \nabla v^2)[a^4(\epsilon + p)] + v(a^4 p)' + \hat{\mathbf{v}} \cdot \nabla(a^4 p) &= 0\end{aligned}$$

conformal time

reduces to **special relativistic Hydrodynamics** when using **rescaled** quantities

How to Calculate Gravitational Waves Analytically?

$$ds^2 = -dt^2 + a^2(\delta_{ij} + h_{ij}(\mathbf{x}))d\mathbf{x}^2$$

Tensor Mode

$$\langle \dot{h}_{ij}(t, \mathbf{q}) \dot{h}_{ij}(t, \mathbf{k}) \rangle = (2\pi)^{-3} \delta^3(\mathbf{k} + \mathbf{q}) P_h(k, t)$$

$$\frac{d\rho_{\text{GW}}(t)}{d \ln k} = \frac{1}{64\pi^3 G} k^3 P_h(t, k)$$

GW Spectrum

Einstein equation

$$h_q'' + 2\frac{a'}{a}h_q' + q^2 h_q = 16\pi G a^2 \pi_q^T$$

neglect backreaction
solve with Green's function

Source evolutions

Plasma(relativistic species), Matter(non-relativistic), Scalar field, EM
Energy-momentum conservation (hydrodynamic limit)

The flow of calculations

$$T_{ij} = a^2 [p\delta_{ij} + (p + e)\gamma^2 v^i v^j]$$



$$\langle \pi_{ij}^T(\eta_1, \mathbf{k}) \pi_{ij}^T(\eta_2, \mathbf{q}) \rangle$$



$$\pi_{ij}^T(\mathbf{q}, \eta) = \frac{a_s^4}{a^4(\eta)} \tilde{\pi}_{ij}^T(\mathbf{q}, \eta)$$



$$\langle \tilde{v}_{\mathbf{q}}^i(\eta_1) \tilde{v}_{\mathbf{k}}^{j*}(\eta_2) \rangle = \delta^3(\mathbf{q} - \mathbf{k}) \hat{q}^i \hat{k}^j G(q, \eta_1, \eta_2)$$

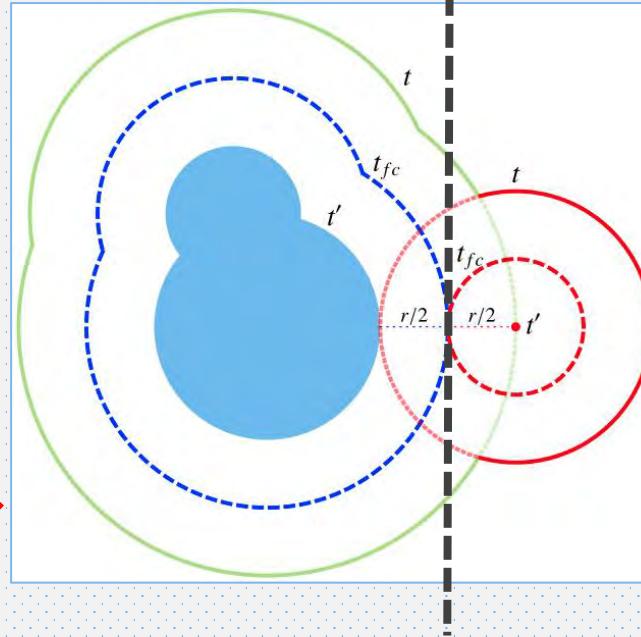
the key part

The Sound Shell Model

- The velocity field is a linear superposition of the contributions from all the bubbles

Hindmarsh, PRL,120,071301,2018, Hindmarsh,Hijazi,JCAP,12,062,2019

contribution from the red bubble



before collision: velocity profile

after collision: sound waves

$$v^i(\eta < \eta_{fc}, \mathbf{x}) = \frac{1}{2} \int \frac{d^3 q}{(2\pi)^3} [\tilde{v}_{\mathbf{q}}^i(\eta) e^{i\mathbf{q} \cdot \mathbf{x}} + \tilde{v}_{\mathbf{q}}^{i*}(\eta) e^{-i\mathbf{q} \cdot \mathbf{x}}]$$

$$v^i(\eta, \mathbf{x}) = \int \frac{d^3 q}{(2\pi)^3} [v_{\mathbf{q}}^i e^{-i\omega\eta + i\mathbf{q} \cdot \mathbf{x}} + v_{\mathbf{q}}^{i*} e^{i\omega\eta - i\mathbf{q} \cdot \mathbf{x}}]$$

$$v_{\mathbf{q}}^i = \sum_{n=1}^{N_b} v_{\mathbf{q}}^{i(n)}$$

Velocity Profile Around a Single Bubble

- Equations in the bag equation of state model can be reduced to Minkowski form
vacuum energy cancels out, junction condition at the wall maintains the same form
- Velocity profile remain unchanged (time \rightarrow conformal time)

Energy fluctuation:

$$(\xi - v)\partial_\xi \tilde{e} = \tilde{w} \left[2\frac{v}{\xi} + \gamma^2(1 - \xi v)\partial_\xi v \right]$$

Pressure fluctuation:

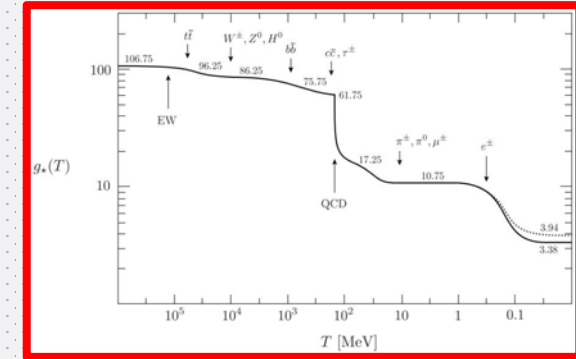
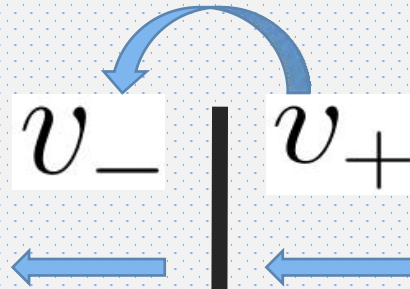
$$(1 - v\xi)\partial_\xi \tilde{p} = \tilde{w}\gamma^2(\xi - v)\partial_\xi v,$$

Velocity equation:

$$2\frac{v}{\xi} = \gamma^2(1 - v\xi) \left[\frac{\mu^2}{c_s^2} - 1 \right] \partial_\xi v$$

self-similar coordinate: $r/\Delta\eta$

Velocity Profile Around a Single Bubble



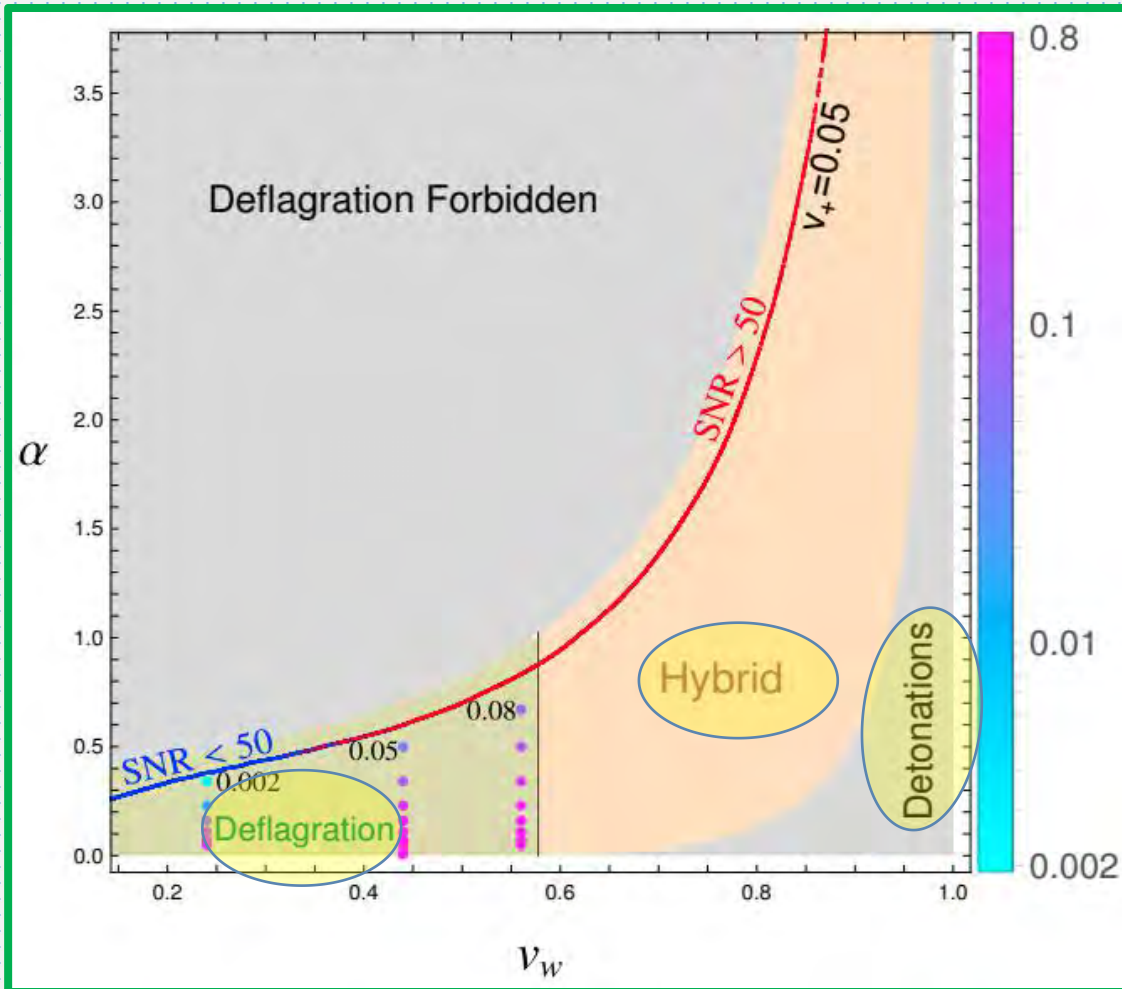
$$p_- = \frac{1}{3}a_-T_-^4, \quad e_- = a_-T_-^4$$

$$p_+ = \frac{1}{3}a_+T_+^4 - \epsilon, \quad e_+ = a_+T_+^4 + \epsilon$$

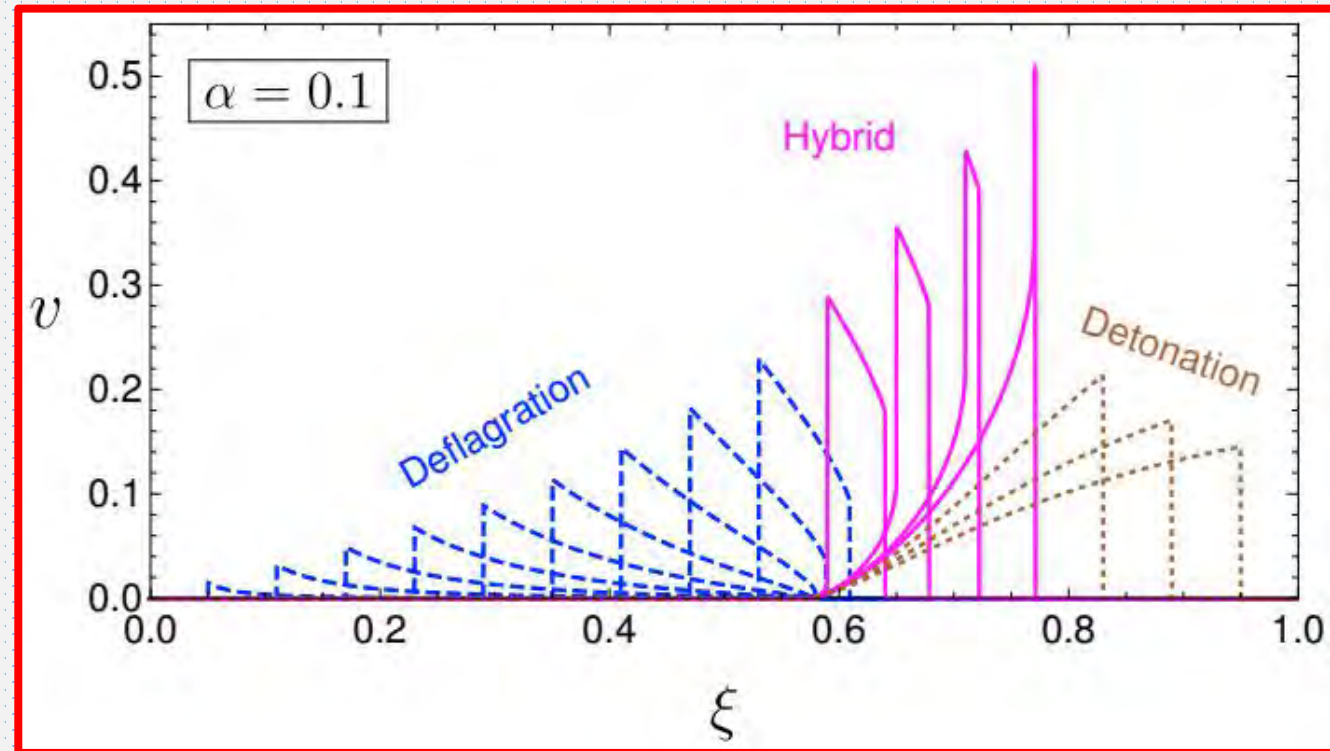
Wall Frame

Velocity Profile Around a Single Bubble

3 Fluid Modes



Velocity Profiles



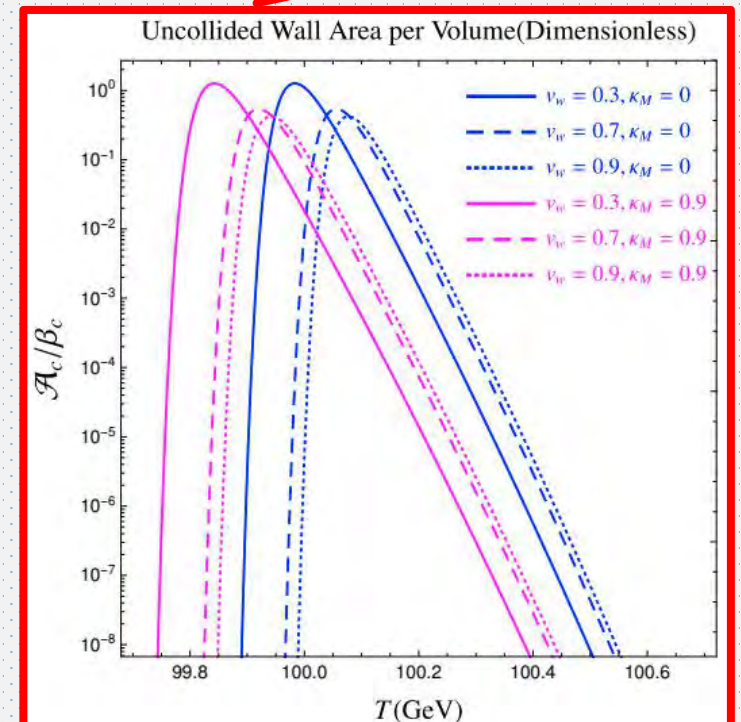
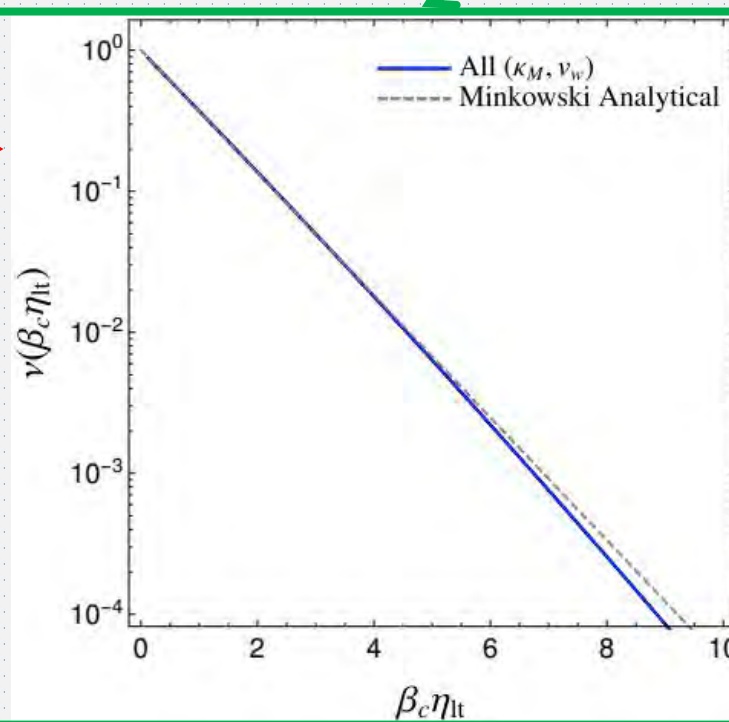
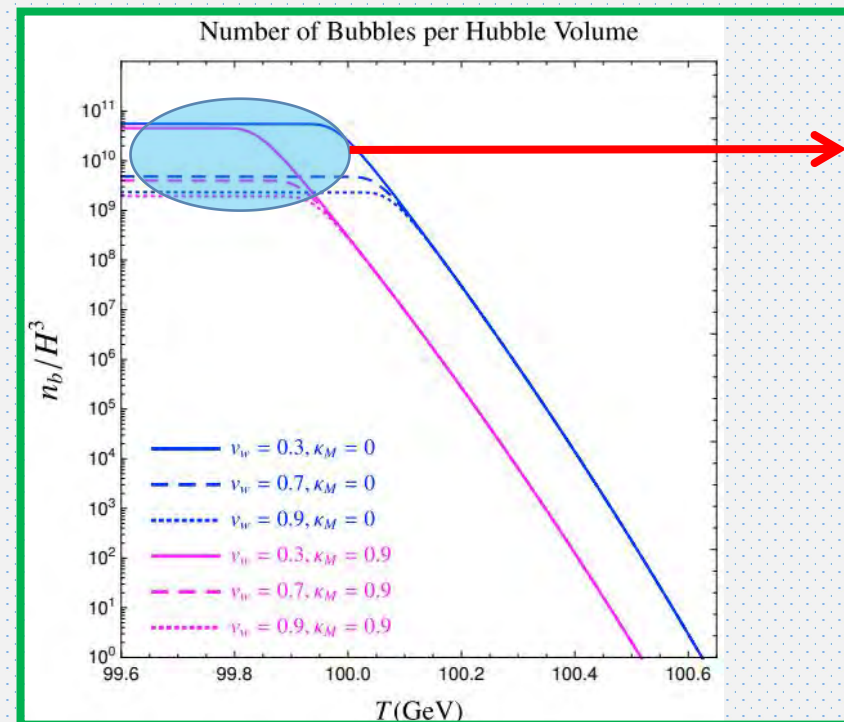
Alves, Ghost, Guo, Sinha, Vagie, JHEP04(2019)052

Velocity Field Power Spectrum

stochastic field: bubble position, formation time, collision time(final size)

after averaging, depends only on conformal lifetime distribution of the bubbles

$$\langle v_{\mathbf{q}_1}^i v_{\mathbf{q}_2}^{j*} \rangle = \hat{q}_1^i \hat{q}_2^j (2\pi)^3 \delta^3(\mathbf{q}_1 - \mathbf{q}_2) \underbrace{\frac{1}{R_{*c}^3 \beta_c^6} \int d\tilde{T} \tilde{T}^6 \nu(\tilde{T}) \left| A\left(\frac{q\tilde{T}}{\beta_c}\right) \right|^2}_{\equiv P_v(q)},$$



Velocity Field Power Spectrum

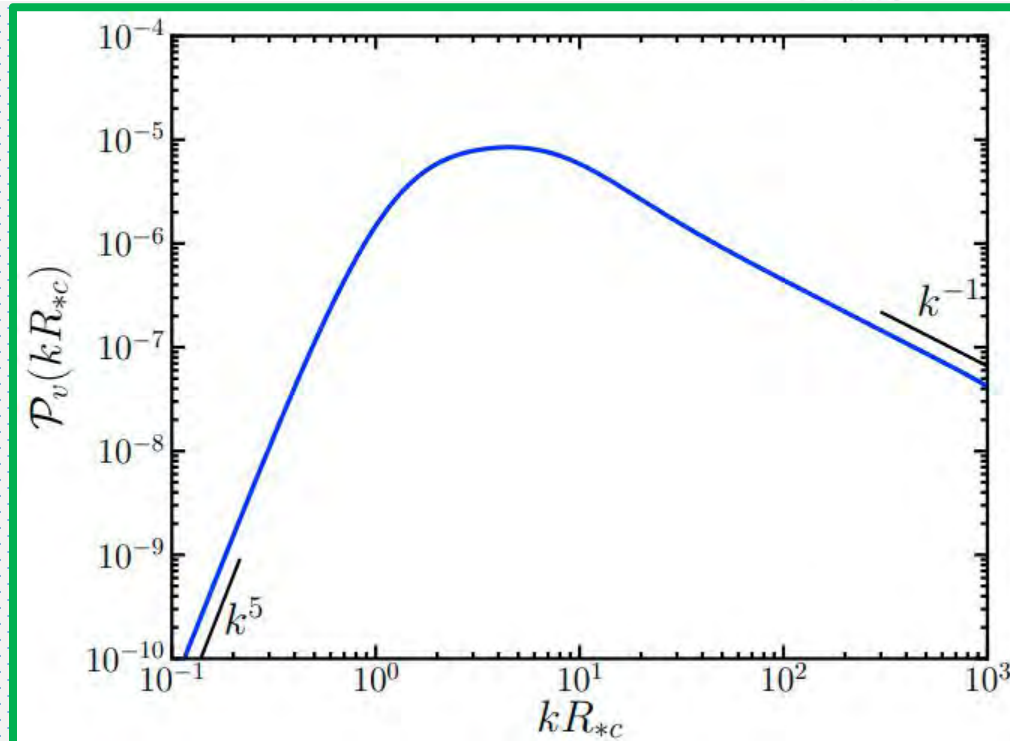
stochastic field: bubble position, formation time, collision time(final size)



after averaging, depends only on conformal lifetime distribution of the bubbles



$$\langle v_{\mathbf{q}_1}^i v_{\mathbf{q}_2}^{j*} \rangle = \hat{q}_1^i \hat{q}_2^j (2\pi)^3 \delta^3(\mathbf{q}_1 - \mathbf{q}_2) \underbrace{\frac{1}{R_{*c}^3 \beta_c^6} \int d\tilde{T} \tilde{T}^6 \nu(\tilde{T}) |A(\frac{q\tilde{T}}{\beta_c})|^2}_{\equiv P_v(q)},$$



Gravitational Wave Power Spectrum

$$(\kappa_M y + 1 - \kappa_M) \frac{d^2 h_q}{dy^2} + \left[\frac{5}{2} \kappa_M + \frac{2(1 - \kappa_M)}{y} \right] \frac{dh_q}{dy} + \tilde{q}^2 h_q = \frac{16\pi G a(y)^2 \pi_q^T(y)}{(a_s H_s)^2}$$

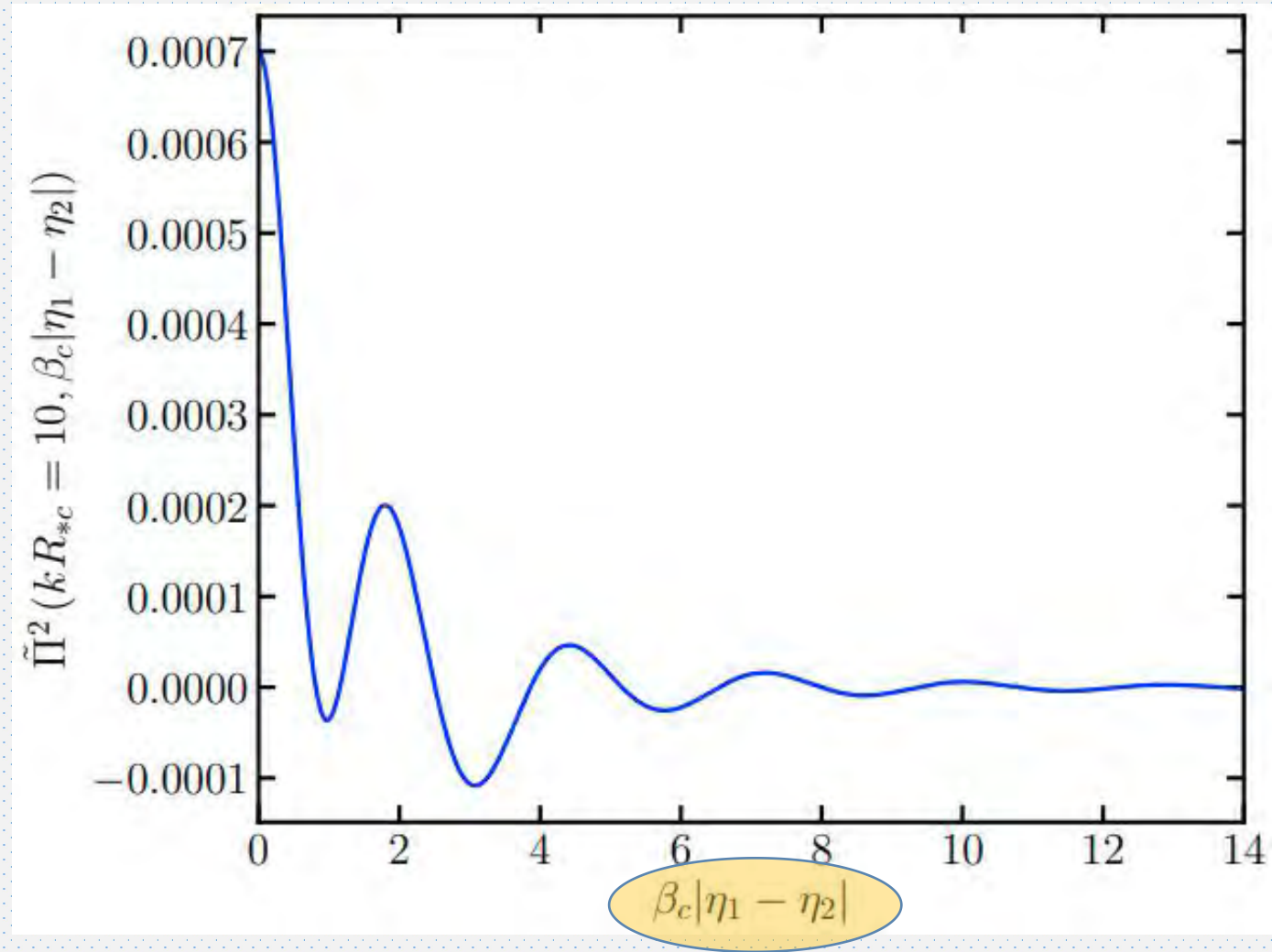
$$y \equiv a/a_s$$

$$h_{ij}(\tilde{y}, \mathbf{q}) = \int_{\tilde{y}_s}^{\tilde{y}} d\tilde{y}' G(\tilde{y}, \tilde{y}') \frac{16\pi G a(\tilde{y}')^2 \pi_{ij}^T(\tilde{y}', \mathbf{q})}{q^2}$$

$$P_{h'} = [16\pi G (\bar{\epsilon} + \bar{p}) \bar{U}_f^2]^2 L_f^3 \int_{\tilde{y}_s}^{\tilde{y}} d\tilde{y}_1 \int_{\tilde{y}_s}^{\tilde{y}} d\tilde{y}_2 \left(\frac{\partial \tilde{y}}{\partial \tilde{\eta}} \right)^2 \frac{\partial G(\tilde{y}, \tilde{y}_1)}{\partial \tilde{y}} \frac{\partial G(\tilde{y}, \tilde{y}_2)}{\partial \tilde{y}} \\ \times \frac{a_s^8}{a^2(\tilde{y}_1) a^2(\tilde{y}_2)} \frac{\tilde{\Pi}^2(k L_f, k \eta_1, k \eta_2)}{k^2}.$$

rescaled, dimensionless unequal time correlator

Gravitational Wave Power Spectrum

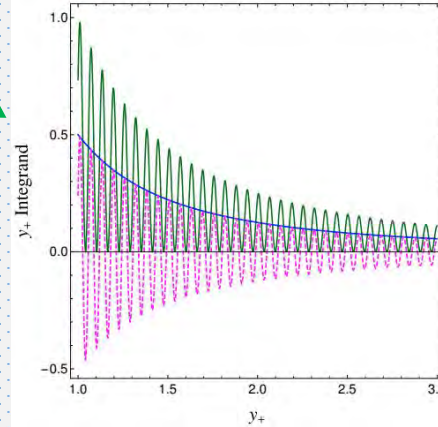


$$y_- \sim \mathcal{O}(10^{-3})/v_w \times \beta_c(\eta_1 - \eta_2) \quad \text{source autocorrelation time very small}$$

Gravitational Wave Power Spectrum

Under reasonable assumptions, the power spectrum can be written down in a similar form as in Minkowski space time

It is also essential to neglect **highly oscillatory part** to obtain the following result



$$\mathcal{P}_{\text{GW}}(y, kR_{*c}) = 3\Gamma^2 \bar{U}_f^4 \frac{H_{R,s}^4}{H^2 H_s} (a_s R_{*c}) \frac{(kR_{*c})^3}{2\pi^2} \tilde{P}_{\text{gw}}(kR_*) \times \frac{1}{y^4} \Upsilon(y)$$

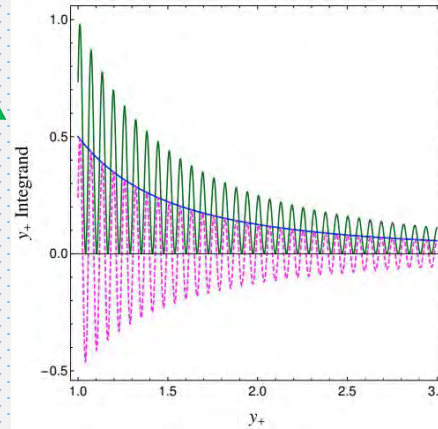
spectrum generally very similar to Minkowski result
but might have additional features when the assumptions can not be made

dilution for modes deep inside the horizon

Gravitational Wave Power Spectrum

Under reasonable assumptions, the power spectrum can be written down in a similar form as in Minkowski space time

It is also essential to neglect **highly oscillatory part** to obtain the following result



effect not captured before

$$\mathcal{P}_{\text{GW}}(y, kR_{*c}) = 3\Gamma^2 \bar{U}_f^4 \frac{H_{R,s}^4}{H^2 H_s} (a_s R_{*c}) \frac{(kR_{*c})^3}{2\pi^2} \tilde{P}_{\text{gw}}(kR_*) \times \frac{1}{y^4} \Upsilon(y)$$

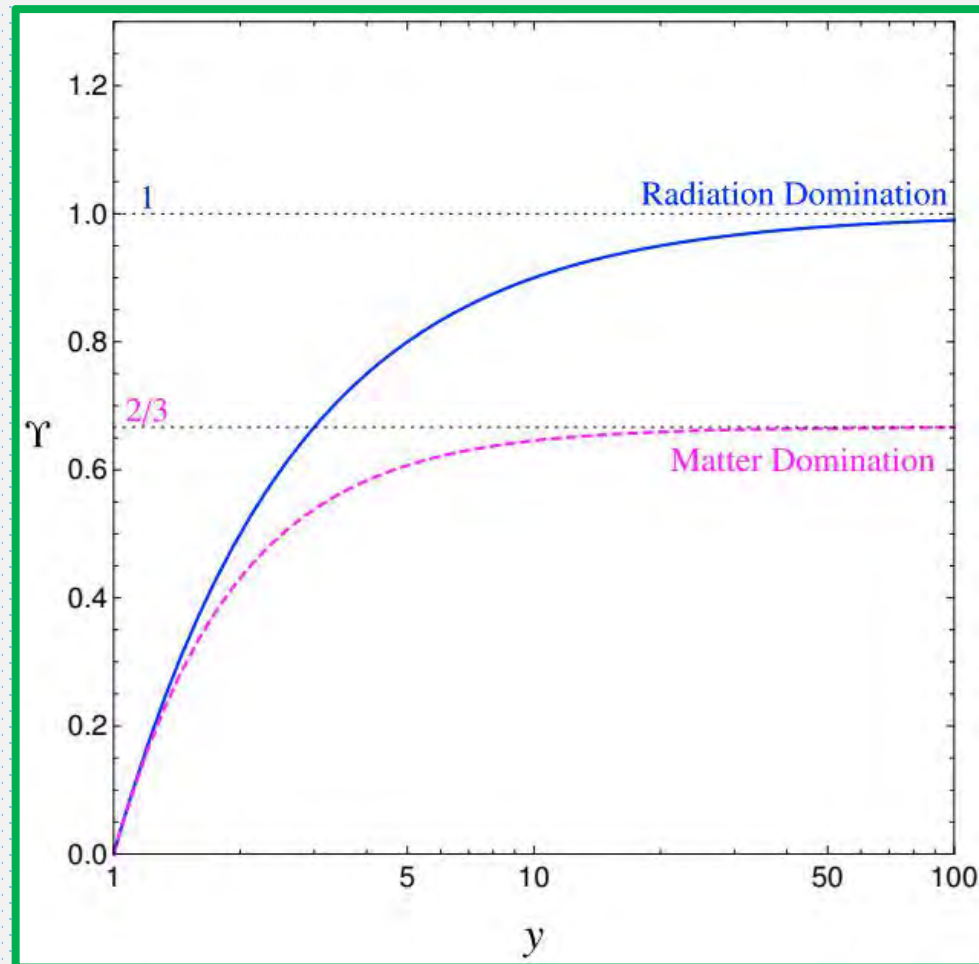
spectrum generally very similar to Minkowski result
but might have additional features when the assumptions can not be made

dilution for modes deep inside the horizon

Gravitational Wave Power Spectrum

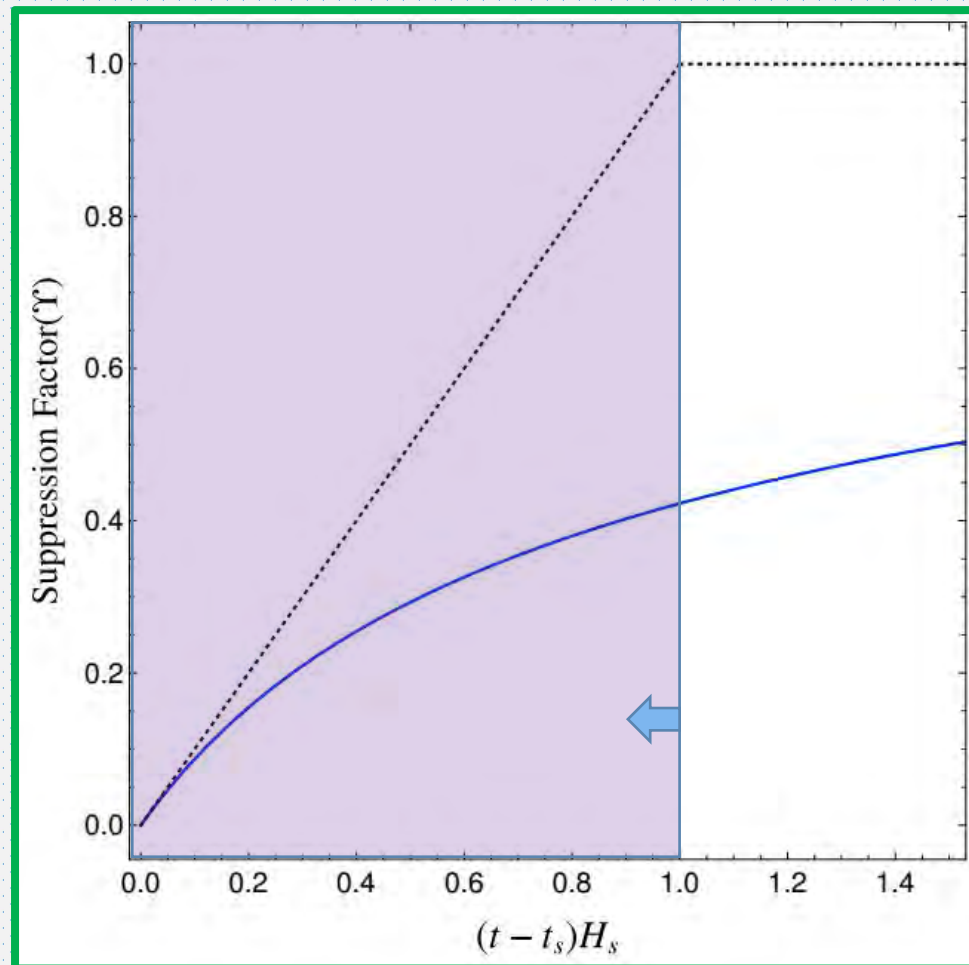
use spectrum from numerical simulation

$$h^2 \Omega_{\text{GW}}(f) = 8.5 \times 10^{-6} \left(\frac{100}{g_s(T_e)} \right)^{1/3} \Gamma^2 \bar{U}_f^4 \left[\frac{H_s}{\beta(v_w)} \right] v_w \mathcal{S}_{\text{SW}}(f) \times \Upsilon(y)$$



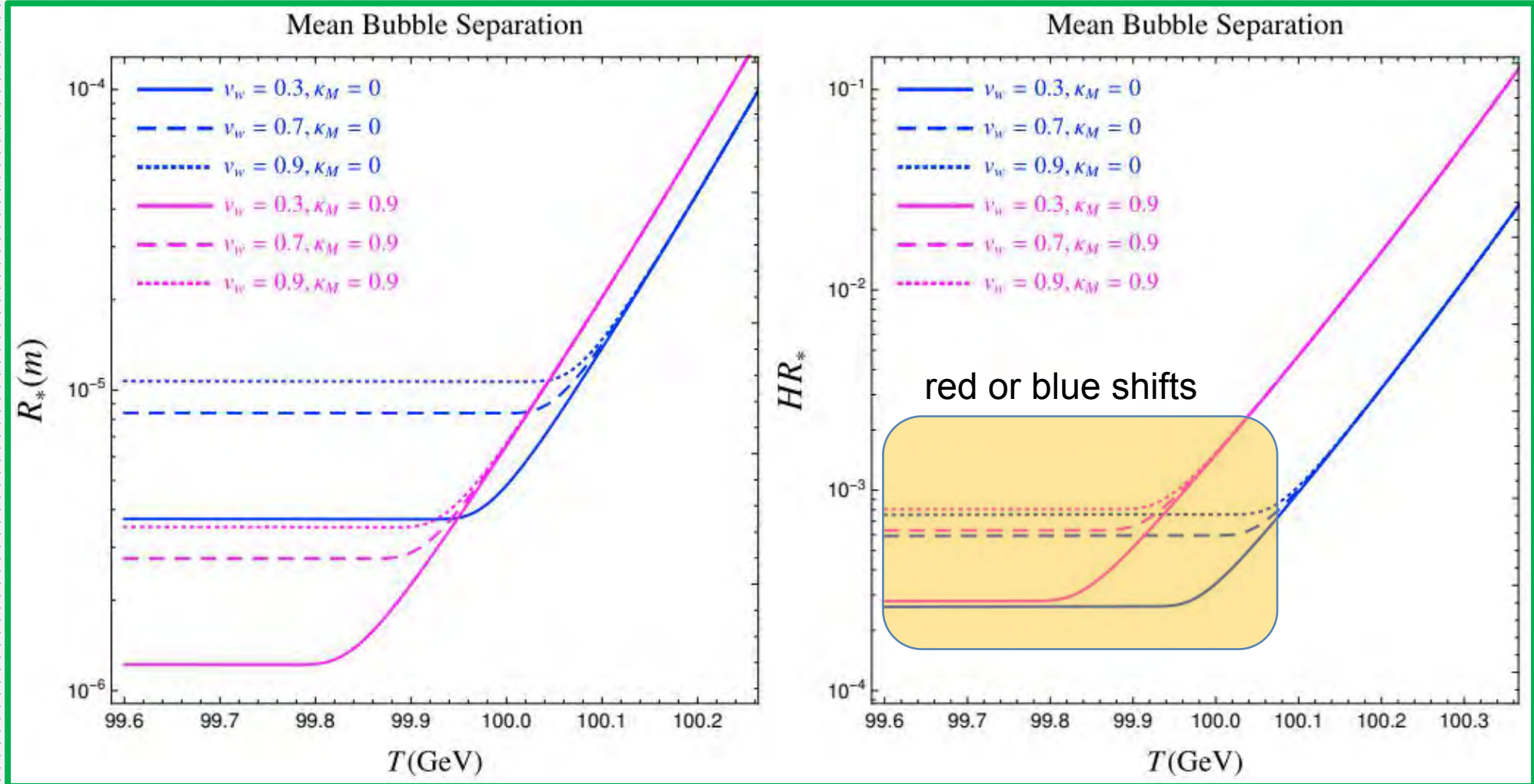
Lifetime of the Source

- Shocks, turbulence, dissipative processes all disrupt the source
- So lifetime is usually less than a Hubble time, meaning a suppression



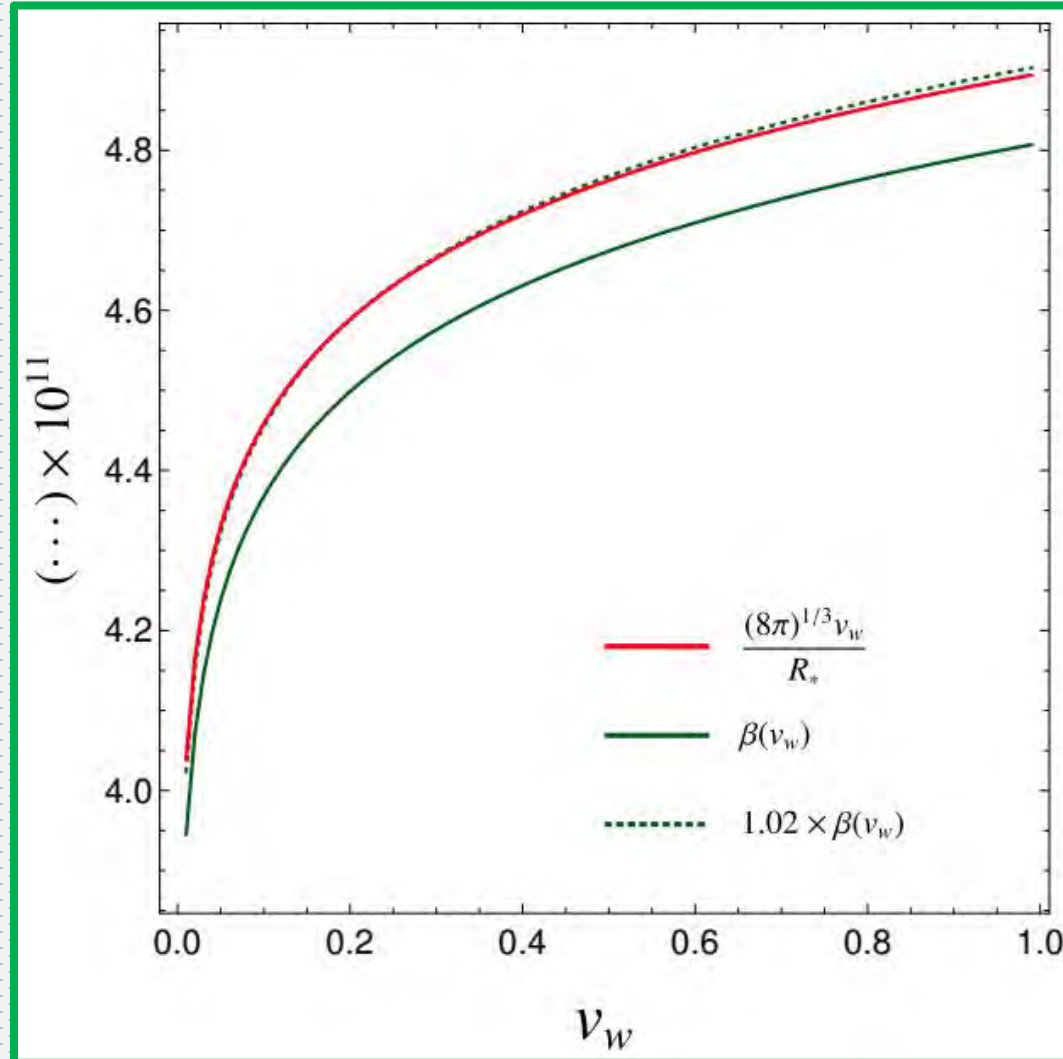
Mean Bubble Separation

The (larger) scale that determines the spectral shape. Peak frequency at $kR_* = z_p$

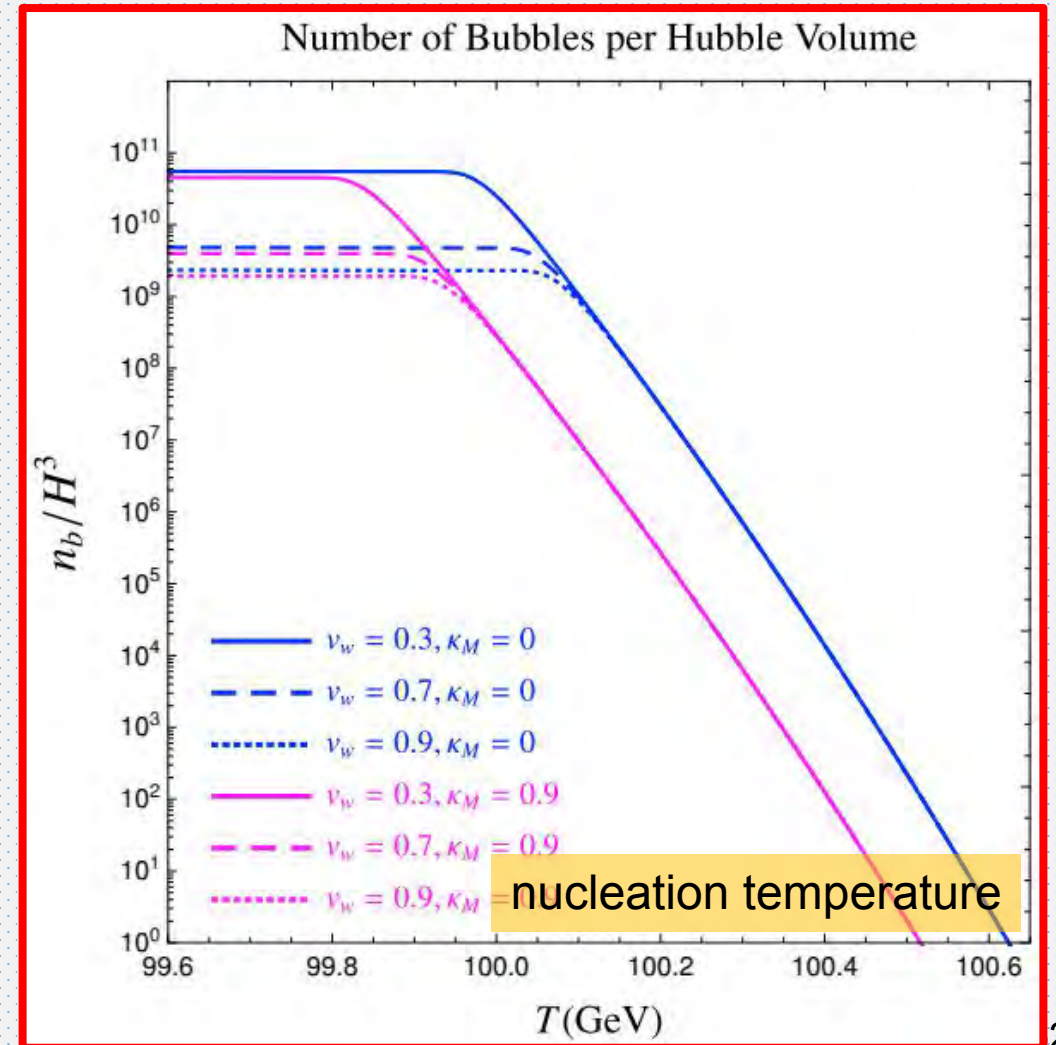
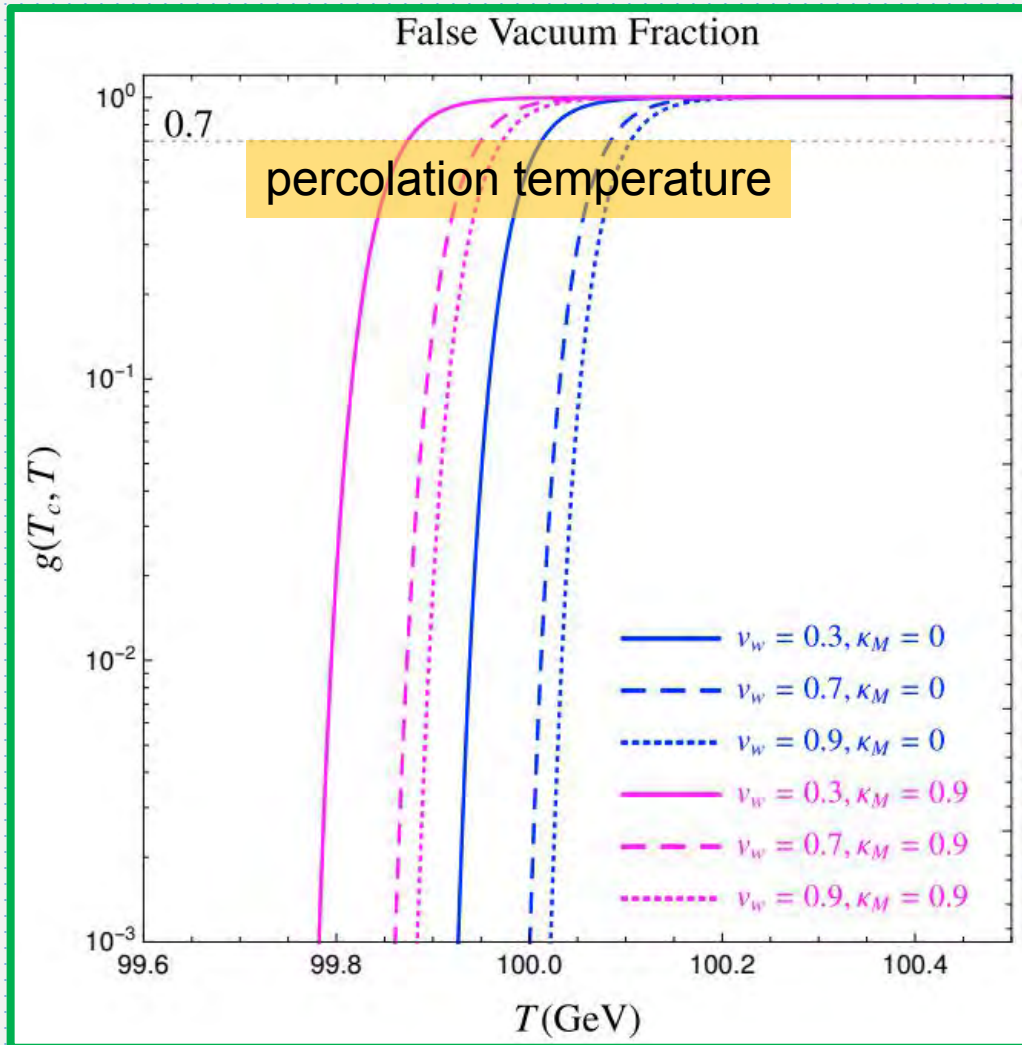


Relation between R^* and Beta

Relation unchanged for our redefined conformal quantities. For physical quantities: $R_*(\eta) = \frac{a(\eta)}{a(\eta_f)} (8\pi)^{1/3} \frac{v_w}{\beta(v_w)}$



Several Temperatures



Summary

- We have set up the framework for modelling the GW from sound waves, in an expanding universe
generally no need for new simulations, rescaled quantities need to be used
- An additional multiplication factor is found and needs to be included to the generally used spectrum
- Details of the PT process is analyzed in standard and non-standard cosmic histories

Thanks!